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UNIFORM l¹ BEHAVIOR OF A TIME DISCRETIZATION METHOD FOR A VOLTERRA INTEGRODIFFERENTIAL EQUATION WITH CONVEX KERNEL; STABILITY*

CHARLES B. HARRIS† AND RICHARD D. NOREN†

Abstract. We study stability of a numerical method in which the backward Euler method is combined with order one convolution quadrature for approximating the integral term of the linear Volterra integrodifferential equation $\mathbf{u}'(t) + \int_0^t \beta(t-s)\mathbf{A}\mathbf{u}(s)\,ds = 0,\ t \geq 0,\ \mathbf{u}(0) = \mathbf{u}_0$, which arises in the theory of linear viscoelasticity. Here \mathbf{A} is a positive self-adjoint densely defined linear operator in a real Hilbert space, and $\beta(t)$ is locally integrable, nonnegative, nonincreasing, convex, and $-\beta'(t)$ is convex. We establish stability of the method under these hypotheses on $\beta(t)$. Thus, the method is stable for a wider class of kernel functions $\beta(t)$ than was previously known. We also extend the class of operators \mathbf{A} for which the method is stable.

Key words. Volterra integrodifferential equation, convolution quadrature, convex kernel, l^1 -stability

AMS subject classifications. 45D05, 45K05, 65R20, 64D05

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1. Introduction. Let **A** be a positive self-adjoint linear operator defined on a dense subspace $\mathcal{D}(\mathbf{A})$ of a real Hilbert space **H** with spectral decomposition

(1)
$$\mathbf{A}\mathbf{x} = \int_{-\infty}^{\infty} \lambda \, d\mathbf{E}_{\lambda} \, \mathbf{x}$$

for $\mathbf{x} \in \mathcal{D}(\mathbf{A})$. We assume that the spectrum of \mathbf{A} is contained in $[\lambda_0, \infty)$, where $\lambda_0 > 0$. Xu established stability results in 2002 (see [21]) and convergence results in 2008 (see [22]) for a numerical method for approximating the initial value problem

(2)
$$\mathbf{u}'(t) + \int_0^t \beta(t-s)\mathbf{A}\mathbf{u}(s) ds = 0, \qquad t \ge 0, \qquad \mathbf{u}(0) = \mathbf{u}_0.$$

Here $\mathbf{u} = \mathbf{u}(t)$ is a function in the Hilbert space \mathbf{H} and ' = d/dt. Xu assumes in both papers that the kernel $\beta(t): (0, \infty) \to \mathbb{R}$ satisfies

(3)
$$\beta \in C(0,\infty) \cap L^1(0,1) \text{ and } 0 \le \beta(\infty) < \beta(0+) \le \infty,$$

and

(4)
$$(-1)^n \beta^{(n)}(t) \ge 0, \quad t > 0, \quad n = 0, 1, 2, \dots$$

In Theorems 1 and 2 we substantially enlarge the class of functions $\beta(t)$ for which the stability results are valid by weakening the completely monotone hypotheses (4) on $\beta(t)$ to the assumption

(5) β is nonnegative, nonincreasing, convex, and $-\beta'$ is convex.

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We also note that our results hold for the wider class of operators **A** defined via a spectral family $\{\mathbf{E}_{\lambda}\}$, as in (1), whereas [21] employed the more restrictive condition that **A** possess a countable complete eigensystem.

Xu utilized a discrete analogue of the Payley-Wiener theorem in [23] to obtain results similar to those in the present paper for a class of quadratures and for certain kernels displaying log convexity. Although the hypotheses in [23] overlap ours, our results hold for kernels lacking log convexity, such as if $\beta(t) = 0$ for some t > 0. As an example,

$$f(x) = \begin{cases} (x_0 - t)^2 & \text{for } 0 \le x \le x_0, \\ 0 & \text{for } x_0 < x \end{cases}$$

for any fixed $x_0 > 0$.

Denote the Laplace transform of a function f by $\widehat{f}(t)$. Thus,

(6)
$$\widehat{\beta}(t) = \int_0^\infty e^{-ts} \beta(s) \, ds, \qquad t > 0.$$

By Bernstein's theorem [20, Chapter 8], a function a = a(t) is completely monotonic iff there exists an associated nonnegative, increasing function $\alpha : [0, \infty) \to [0, \infty)$ with

(7)
$$a(t) = \int_0^\infty e^{-xt} d\alpha(x), \qquad t > 0.$$

From (7) we see that the Laplace transform of a(t) may be analytically extended to the slit plane $\mathbb{C}' \equiv \mathbb{C} \setminus (-\infty, 0]$ via the formula

(8)
$$\widehat{a}(t) = \int_0^\infty \frac{d\alpha(s)}{s+t} \qquad t \in \mathbb{C}'.$$

Here a Stieltjes integral is used. Xu makes extensive use of this representation in his analysis.

A convex function will only be guaranteed to have a second derivative almost everywhere [18, Chapter 7]. In particular, the representation (8) does not hold. Without this representation we are still able to obtain the same conclusions as Xu by doing detailed estimates on the function $\hat{\beta}(t)$ using the representation (6).

Let k denote the constant time step, $t_n = kn$ the nth time level, and \mathbf{U}^n the approximation of $\mathbf{u}(t_n)$. The backward Euler method is used with $\bar{\partial} \mathbf{U}^n = \frac{\mathbf{U}^n - \mathbf{U}^{n-1}}{k}$ approximating the derivative \mathbf{u}' in (2) at the nth time level. For the integral we apply the first-order convolution quadrature introduced by Lubich [7]:

(9)
$$q_n(\varphi) = \sum_{i=1}^n \beta_{n-j}(k) \varphi^j,$$

where $\varphi^j = \varphi(t_j)$ and the quadrature weights $\beta_{n-j}(k)$ are the coefficients of the power series

(10)
$$\widehat{\beta}\left(\frac{1-z}{k}\right) = \sum_{j=0}^{\infty} \beta_j(k)z^j.$$

This leads to the time discrete problem

(11)
$$\bar{\partial} \mathbf{U}^n + q_n(\mathbf{A}\mathbf{U}) = 0, \qquad \mathbf{U}^0 = \mathbf{u}_0.$$

Our first theorem generalizes Theorem 1 of [21] by replacing the completely monotonic assumption (4) with (5).

THEOREM 1. If (3) and (5) hold, then

(12)
$$k \sum_{n=1}^{\infty} ||\mathbf{U}^n|| \le C||\mathbf{A}\mathbf{u}_0||.$$

In order to state our next theorem we must first define some auxiliary functions. For $\sigma + i\tau \notin (-\infty, 0]$, set $\beta(t) = c(t) + \beta(\infty)$, and then let

$$\phi(\sigma,\tau) = \int_0^\infty \cos(\tau t)e^{-\sigma t}\beta(t) dt \quad \text{and} \quad \theta(\sigma,\tau) = \frac{1}{\tau} \int_0^\infty \sin(\tau t)e^{-\sigma t}\beta(t) dt,$$

$$\phi_c(\sigma,\tau) = \int_0^\infty \cos(\tau t)e^{-\sigma t}c(t) dt \quad \text{and} \quad \theta_c(\sigma,\tau) = \frac{1}{\tau} \int_0^\infty \sin(\tau t)e^{-\sigma t}c(t) dt,$$

and for $0 < \tau < \infty$, set

$$\phi_c(\tau) = \lim_{\sigma \to 0+} \phi_c(\sigma, \tau) = \int_0^\infty \cos(\tau t) c(t) dt$$

and

$$\theta_c(\tau) = \lim_{\sigma \to 0+} \theta_c(\sigma, \tau) = \frac{1}{\tau} \int_0^\infty \sin(\tau t) c(t) dt.$$

So, for $\sigma + i\tau \notin (-\infty, 0]$, we have

$$\phi(\sigma,\tau) = \phi_c(\sigma,\tau) + \frac{\sigma\beta(\infty)}{\sigma^2 + \tau^2}$$
 and $\theta(\sigma,\tau) = \theta_c(\sigma,\tau) + \frac{\beta(\infty)}{\sigma^2 + \tau^2}$,

and then, for $0 < \tau < \infty$, we may set

$$\phi(\tau) = \lim_{\sigma \to 0+} \phi(\sigma, \tau) = \phi_c(\tau)$$
 and $\theta(\tau) = \lim_{\sigma \to 0+} \theta(\sigma, \tau) = \theta_c(\tau) + \frac{\beta(\infty)}{\tau^2}$

We see then that the Fourier transform of $\beta(t)$,

(13)
$$\widetilde{\beta}(\tau) = \int_0^\infty e^{-i\tau t} \beta(t) dt,$$

obeys the relation

(14)
$$\widetilde{\beta}(\tau) = \phi(\tau) - i\tau\theta(\tau),$$

and further, the Laplace transform obeys

(15)
$$\widehat{\beta}(\sigma + i\tau) = \phi(\sigma, \tau) - i\tau\theta(\sigma, \tau).$$

As a consequence of Theorem 2.2 and Corollary 2.1 of Carr and Hannsgen [1], (3) and (5) imply

(16)
$$\limsup_{\tau \to \infty} \frac{\theta_c(\tau)}{\phi_c(\tau)} < \infty.$$

By (4.3) of [1], we see that $\tau^{-2} = o(\theta_c(\tau))$ $(\tau \to \infty)$, so it follows that (3) and (5) imply that

(17)
$$\limsup_{\tau \to \infty} \frac{\theta(\tau)}{\phi(\tau)} < \infty.$$

If instead our kernel $\beta(t)$ is such that

(18)
$$\limsup_{\tau \to \infty} \frac{\tau^{\frac{1}{3}}\theta(\tau)}{\phi(\tau)} < \infty$$

holds, then we can obtain the following theorem which generalizes Theorem 2 of [21]. THEOREM 2. If (3), (5), and (18) hold, then

(19)
$$k\sum_{n=1}^{\infty}||\mathbf{U}^n|| \le C||\mathbf{u}_0||.$$

We note that (18) is a significantly weaker frequency condition upon $\beta(t)$ than is employed in Theorem 2 of [21], namely, that

(20)
$$\limsup_{\tau \to \infty} \frac{\tau \theta(\tau)}{\phi(\tau)} < \infty.$$

For example, if $\beta(t)$ satisfies (5) and behaves like $(-\log(t))^p$ (p > 0) near the origin, then an easy calculation utilizing the relations (37) and (39) shows that (18) is satisfied, but not (20). We see that in Theorem 1 we are allowed a wider class of kernel functions $\beta(t)$, but we have the more restrictive requirement that $\mathbf{u}_0 \in \mathcal{D}(\mathbf{A})$, whereas Theorem 2 places greater restrictions upon $\beta(t)$, yet allows \mathbf{u}_0 to be any element of \mathbf{H} .

The resolvent kernel of (2) is defined by the formula

(21)
$$\mathbf{U}(t) = \int_{-\infty}^{\infty} u(t,\lambda) d\mathbf{E}_{\lambda},$$

where $u(t,\lambda)$ is the solution of the scalar Volterra integrodifferential equation

(22)
$$u'(t) + \lambda \int_0^t \beta(t-s)u(s) \, ds = 0, \quad u(0) = 1;$$

the parameter λ satisfies $\lambda_0 \leq \lambda$ and $0 \leq t$.

It is clear from (21) that

(23)
$$\sup_{\lambda_0 \le \lambda} |u(t,\lambda)| \to 0, \quad t \to \infty$$

and

(24)
$$\int_{0}^{\infty} \sup_{\lambda_0 \le \lambda} |u(t,\lambda)| \, dt < \infty$$

imply, respectively,

(25)
$$\|\mathbf{U}(t)\| \to 0, \ t \to \infty$$

and

(26)
$$\int_0^\infty \|\mathbf{U}(t)\| \, dt < \infty.$$

Then the resolvent formula

(27)
$$\mathbf{y}(t) = \mathbf{U}(t)\mathbf{y}_0 + \int_0^t \mathbf{U}(t-s)\mathbf{f}(s) ds$$

can be used to obtain precise asymptotic information $(t \to \infty)$ about the solution $\mathbf{y}(t)$ of the initial value problem

(28)
$$\mathbf{u}'(t) + \int_0^t \beta(t-s)\mathbf{A}\mathbf{u}(s) ds = \mathbf{f}(t), \qquad t \ge 0, \qquad \mathbf{u}(0) = \mathbf{u}_0.$$

In [1] several sufficient conditions are given on $\beta(t)$ such that (23) and (24) hold. One easily stated consequence of [1] which is relevant here is that (23) and (24) both hold, and, as a consequence, (25) and (26) when $\beta(t)$ satisfies (5).

In [21] the stability of a numerical scheme for approximating the solution of (2) is a discrete analogue of (26). Let $\{U^n(\lambda)\}_{n=0}^{\infty}$ be a real sequence satisfying the difference equation

(29)
$$\frac{U^n(\lambda) - U^{n-1}(\lambda)}{k} + \lambda q_n(U(\lambda)) = 0, \qquad n \ge 1, \qquad U^0(\lambda) = 1.$$

It follows from the functional calculus for spectral decompositions (see [17]) that the solution to (11) may be representated as

(30)
$$\mathbf{U}^n = \int_{-\infty}^{\infty} U^n(\lambda) \, d\mathbf{E}_{\lambda} \, \mathbf{u}_0.$$

We note that Lemma 1 from [6] implies that $e^{-\sigma t}\beta(t)$ and $(e^{-\sigma t}\beta(t))'$ are convex for $\sigma > 0$. Also, from Theorem 2 and the comments following it in [13] we find that $\beta(t)$ is positive-definite, implying that $\operatorname{Re}(\widehat{\beta}(s)) > 0$ whenever $s = \sigma + i\tau$ with $\sigma > 0$. Then, by an argument similar to that in Lemma 3.1 of [8], we find that the quadrature (9) is positive-definite in the sense that for each function $\varphi : (0, \infty) \to \mathbf{H}$ and each positive integer N, we have

(31)
$$Q_N(\varphi) \equiv k \sum_{n=1}^N (q_n(\varphi), \varphi^n) \ge 0.$$

To see this, set

$$\widetilde{\varphi}(t) = \sum_{j=1}^{N} \varphi^{j} t^{j}, \quad \widetilde{\beta}(t) = \sum_{j=0}^{\infty} \beta_{j}(k) t^{j} \quad \text{and} \quad Q_{N,r}(\varphi) = k \sum_{n=1}^{N} \sum_{j=1}^{n} \beta_{n-j}(k) r^{n-j} (\varphi^{j}, \varphi^{n})$$

for 0 < r < 1. Then, it is straightforward to show that

$$Q_{N,r}(\varphi) = \frac{k}{2\pi} \int_0^{2\pi} \widetilde{\beta}(re^{i\theta}) ||\widetilde{\varphi}(e^{i\theta})||^2 d\theta.$$

As **H** is a real Hilbert space, it follows from (10) that $Q_{N,r}(\varphi) \geq 0$. Then, by (9) we find that $Q_{N,r}(\varphi) \to Q_N(\varphi)$ $(r \uparrow 1)$, from which (31) follows.

By an argument very similar to that given in Lemma 3.1 of [10], it can be shown that (31) implies that

$$(32) ||\mathbf{U}^n|| \le ||\mathbf{u}_0||.$$

Then (32) implies that

(33)
$$k \sum_{n=1}^{m} ||\mathbf{U}^n|| \le t_m ||\mathbf{u}_0||.$$

Thus, by (30) and (33), we see that it is sufficient to show that

(34)
$$k \sum_{n=m+1}^{\infty} \sup_{\lambda \ge \lambda_0} |U^n(\lambda)\lambda^{-1}| \le C$$

and

(35)
$$k \sum_{n=m+1}^{\infty} \sup_{\lambda \ge \lambda_0} |U^n(\lambda)| \le C$$

to prove Theorems 1 and 2, respectively.

Equations (2) and (28) arise in the theory of linear viscoelasticity. A nice survey may be found in [16]. For a comprehensive treatment of Volterra equations see [5] or [15]. Another interesting work on the numerical approximation of the solution of (2) which assumes (3) and (5) is given by Xu in [24, Remark 2.3] in which a Galerkin method is studied. For a numerical solution utilizing quadrature applied to the inverse Laplace transform form of the solution, see [11]. For a second-order accurate finite difference solution, see [9]. A solution utilizing finite difference convolution quadrature is given in [3]. For a time-stepping discontinuous Galerkin solution, see [12].

In the next section we establish some preliminary results and in section 3 we present the proofs of our theorems. In all that follows we assume that $\varepsilon > 0$ is a sufficiently small fixed constant independent of k whose value will be specified later. We also note that C is a generic constant whose value may change at each appearance and which depends only upon ε and λ_0 .

2. Preliminary estimates. We begin with a lemma from [21, p. 139], which derives from a lemma in [1, p. 967].

LEMMA 2.1. If $\beta(t)$ satisfies (3) and (5), then

(36)
$$\frac{1}{2\sqrt{2}} \int_0^{\frac{1}{\tau}} \beta(t) \, dt \le |\widetilde{\beta}(\tau)| \le 4 \int_0^{\frac{1}{\tau}} \beta(t) \, dt, \quad \tau > 0,$$

(37)
$$\frac{1}{5} \int_0^{\frac{1}{\tau}} t\beta(t) dt \le \theta(\tau) \le 12 \int_0^{\frac{1}{\tau}} t\beta(t) dt, \quad \tau > 0,$$

(38)
$$|\widetilde{\beta}'(\tau)| \le 40 \int_0^{\frac{1}{\tau}} t\beta(t) dt, \quad \tau > 0.$$

Here, recall that $\widetilde{\beta}(\tau)$ is the Fourier transform of $\beta(t)$. We note that these results hold without the convexity of $-\beta'(t)$ assumed. As we know that $e^{-\sigma t}\beta(t)$ and $(e^{-\sigma t}\beta(t))'$ are convex for $\sigma > 0$, then with only slight modifications to the proof we obtain results similar to those in Noren (see [14, eq. (4.14)]):

(39)
$$\frac{1}{C} \int_0^{\frac{1}{\tau}} -t\beta'(t) dt \le \phi(\tau) \le C \int_0^{\frac{1}{\tau}} -t\beta'(t) dt, \quad \tau > 0,$$

and

(40)
$$\frac{1}{C} \int_{0}^{\frac{1}{\tau}} -t(e^{-\sigma t}\beta(t))' dt \le \phi(\sigma,\tau) \le C \int_{0}^{\frac{1}{\tau}} -t(e^{-\sigma t}\beta(t))' dt, \quad \sigma,\tau > 0.$$

One consequence of (39) and (40) in the case where $0 < \sigma \le \varepsilon \tau$ is that

(41)
$$\phi(\sigma,\tau) \ge C \int_0^{\frac{1}{\tau}} -t(e^{-\sigma t}\beta(t))' dt \ge Ce^{-\frac{\sigma}{\tau}} \int_0^{\frac{1}{\tau}} -t\beta'(t) dt \ge C\phi(\tau).$$

As $e^{-\sigma t}\beta(t)$ satisfies the hypotheses of Lemma 2.1, we obtain the following variants of (36) and (38):

$$(42) \qquad \frac{1}{2\sqrt{2}} \int_{0}^{\frac{1}{\tau}} e^{-\sigma t} \beta(t) \, dt \le |(e^{-\sigma t} \beta(t))^{\sim}(\tau)| \le 4 \int_{0}^{\frac{1}{\tau}} e^{-\sigma t} \beta(t) \, dt, \qquad \sigma, \tau > 0,$$

and

(43)
$$\left| \frac{d}{d\tau} (e^{-\sigma t} \beta(t))^{\sim} (\tau) \right| \le 40 \int_0^{\frac{1}{\tau}} t e^{-\sigma t} \beta(t) dt, \quad \sigma, \tau > 0.$$

Defining the functions $A(x) = \int_0^x \beta(t) dt$ and $A_1(x) = \int_0^x t \beta(t) dt$, we also recall a result from Shea and Wainger [19, eq. (1.21)]):

(44)
$$\int_0^\varepsilon \frac{A_1(\tau^{-1})}{A^2(\tau^{-1})} d\tau < \infty.$$

Define the notations

$$\sigma = \sigma(k, \nu) = \frac{1 - \cos(k\nu)}{k}, \ \tau = \tau(k, \nu) = \frac{\sin(k\nu)}{k}, \ s = s(k, \nu) = \sigma + i\tau = \frac{1 - e^{-ik\nu}}{k},$$

and

$$D(s,\lambda) = D(\sigma + i\tau) = \frac{s}{\lambda} + \widehat{\beta}(s) = \frac{\sigma + i\tau}{\lambda} + \widehat{\beta}(\sigma + i\tau).$$

Following [1] and [21], we wish to define a strictly increasing function $\omega: [\lambda_0, \infty) \to [\varepsilon, \infty)$ with $\omega(\lambda) \to \infty$ ($\lambda \to \infty$) and such that $\theta(\omega(\lambda)) = \frac{1}{\lambda}$ for $\lambda \ge \lambda_1 \ge \lambda_0$ and, if necessary, $\omega(\lambda) = \varepsilon$ for $\lambda_1 > \lambda \ge \lambda_0$. We note that ω was continuous in [1], owing to the choice of $\rho = \frac{6}{t_1}$ in that paper, and in [21] by the analytic nature of a completely monotonic function. We do not require that ω be continuous. In this case, slight modification to the proof given in [1, Lemma 5.2] and [2, Lemma 8.1] gives us the following lemma.

LEMMA 2.2. If $\beta(t)$ satisfies (3) and (5), then

(45)
$$|D(i\tau,\lambda)| \ge \begin{cases} C_1 \frac{|\tau - \omega|}{\lambda} & (\tau \ge \frac{\omega}{2}), \\ C_1(\tau \int_0^{\frac{1}{\tau}} t\beta(t) dt + \int_0^{\frac{1}{\tau}} \beta(t) dt) & (\tau \in [\frac{\varepsilon}{2}, \frac{\omega}{2}]). \end{cases}$$

This result also holds if $-\beta'(t)$ convex is dropped. We also note that [1] gives us

and it follows from (6.8) of [1] that, for $\tau \geq \frac{\omega}{2}$, we have

(47)
$$\theta(\tau) \le C\lambda^{-1}.$$

We now wish to establish a generalization of (2.9) from [21]. LEMMA 2.3. If $\beta(t)$ satisfies (3) and (5) and $0 < \sigma \le \varepsilon \tau < \tau$, then

(48)
$$|\theta(\sigma,\tau) - \theta(\tau)| \le 29\varepsilon\theta(\tau).$$

Proof. Beginning with the formulas

$$\theta(\sigma, \tau) = \theta_c(\sigma, \tau) + \frac{\beta(\infty)}{\sigma^2 + \tau^2}$$
 and $\theta(\tau) = \theta_c(\tau) + \frac{\beta(\infty)}{\tau^2}$,

we see that

$$|\theta(\sigma,\tau) - \theta(\tau)| \le |\theta_c(\sigma,\tau) - \theta_c(\tau)| + \beta(\infty) \left| \frac{1}{\tau^2} - \frac{1}{\sigma^2 + \tau^2} \right|$$

$$\le |\theta_c(\sigma,\tau) - \theta_c(\tau)| + \beta(\infty) \frac{\varepsilon}{\tau^2},$$

so it suffices for us to show that $|\theta_c(\sigma,\tau) - \theta_c(\tau)| \leq 29\varepsilon\theta_c(\tau)$. Integrating by parts twice, we get

$$\theta_c(\sigma,\tau) = \frac{1}{\sigma^2 + \tau^2} \int_0^\infty \left\{ (1 - e^{-\sigma t} \cos(\tau t)) - \frac{\sigma}{\tau} e^{-\sigma t} \sin(\tau t) \right\} (-c'(t)) dt$$

$$= \frac{1}{(\sigma^2 + \tau^2)^2} \int_0^\infty \left\{ \left((\sigma^2 + \tau^2)t + (\sigma^2 - \tau^2)e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right) - 2\sigma (1 - e^{-\sigma t} \cos(\tau t)) \right\} c''(t) dt$$

and

$$\theta_c(\tau) = \frac{1}{\tau^2} \int_0^\infty (1 - \cos(\tau t))(-c'(t)) dt$$
$$= \frac{1}{\tau^2} \int_0^\infty \left(t - \frac{\sin(\tau t)}{\tau} \right) c''(t) dt.$$

The boundary terms vanish due to the relations $tc(t) = t^2c'(t) = o(1)$ $(t \to 0+)$ and tc'(t) = o(1) $(t \to \infty)$ from [1]. Then, setting

$$f(t) = \frac{1}{(\sigma^2 + \tau^2)^2} \left\{ \left((\sigma^2 + \tau^2)t + (\sigma^2 - \tau^2)e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right) - 2\sigma (1 - e^{-\sigma t} \cos(\tau t)) \right\}$$

and

$$g(t) = \frac{1}{\tau^2} \left(t - \frac{\sin(\tau t)}{\tau} \right),$$

we see that we need only show that

$$(49) (1 - 29\varepsilon)g(t) \le f(t) \le (1 + 29\varepsilon)g(t)$$

to have our result. Since f'(0) = f(0) = g'(0) = g(0) = 0 and $(1 - \varepsilon)g''(t) \le f''(t) \le g''(t)$ for $t \in [0, \frac{1}{\tau}]$, we may integrate twice over [0, t] for $t \in [0, \frac{1}{\tau}]$ to obtain $(1 - \varepsilon)g(t) \le f(t) \le g(t)$ for $t \in [0, \frac{1}{\tau}]$. First we show that

(50)
$$(1 - 29\varepsilon)g(t) \le f(t) \quad \left(t > \frac{1}{\tau}\right).$$

Note first that as $0 < \sigma \le \varepsilon \tau < \tau$ and $t > \frac{1}{\tau}$, we have

$$\frac{-2\sigma}{(\sigma^2+\tau^2)^2}(1-e^{-\sigma t}\cos(\tau t))\geq \frac{-2\sigma}{\tau^4}(1-e^{-\sigma t}\cos(\tau t))\geq \frac{-4\varepsilon}{\tau^3}\geq \frac{-26\varepsilon}{\tau^2}\left(t-\frac{\sin(\tau t)}{\tau}\right).$$

So, we must show that

$$\frac{1}{(\sigma^2 + \tau^2)^2} \left((\sigma^2 + \tau^2)t + (\sigma^2 - \tau^2)e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right) \ge (1 - 3\varepsilon) \left(\frac{1}{\tau^2} \left(t - \frac{\sin(\tau t)}{\tau} \right) \right).$$

As $(1+\varepsilon)(1-3\varepsilon) \le (1-2\varepsilon)$ and

$$\begin{split} \frac{1}{(\sigma^2 + \tau^2)^2} \left((\sigma^2 + \tau^2)t + (\sigma^2 - \tau^2)e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right) \\ & \geq \frac{1}{(1 + \varepsilon)\tau^2} \left(t - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right), \end{split}$$

it suffices to show that (after some rearrangement)

$$\left(1 - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t}\right) \frac{\sin(\tau t)}{\tau} \ge 2\varepsilon \left(\frac{\sin(\tau t)}{\tau} - t\right).$$

This clearly holds if $\sin(\tau t) \geq 0$, so assume otherwise. Then, as

$$\left(1 - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t}\right) \frac{\sin(\tau t)}{\tau} \ge \left(1 - \frac{1 - \varepsilon}{1 + \varepsilon} (1 - \varepsilon \tau t)\right) \frac{\sin(\tau t)}{\tau}
= \left(\frac{2\varepsilon}{1 + \varepsilon} + \left(\frac{1 - \varepsilon}{1 + \varepsilon}\right) \varepsilon \tau t\right) \frac{\sin(\tau t)}{\tau},$$

(50) follows since $2\varepsilon \leq 2\varepsilon(1+\varepsilon)$ and $(\frac{1-\varepsilon}{1+\varepsilon})\sin(\tau t) \geq -1$. We now show that

(51)
$$f(t) \le (1 + 29\varepsilon)g(t) \quad \left(t > \frac{1}{\tau}\right).$$

Note that as $0 < \sigma \le \varepsilon \tau < \tau$ and $t > \frac{1}{\tau}$, we have

$$f(t) \le \frac{1}{\sigma^2 + \tau^2} \left(t - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right) \le \frac{1}{\tau^2} \left(t - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t} \frac{\sin(\tau t)}{\tau} \right).$$

So, we need only show that (after some rearrangement)

$$\left(1 - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t}\right) \frac{\sin(\tau t)}{\tau} \le 29\varepsilon \left(t - \frac{\sin(\tau t)}{\tau}\right).$$

This clearly holds if $\sin(\tau t) \leq 0$, so assume otherwise. Then, we see that

$$\left(1 - \frac{\tau^2 - \sigma^2}{\tau^2 + \sigma^2} e^{-\sigma t}\right) \frac{\sin(\tau t)}{\tau} \le \left(1 - \frac{1 - \varepsilon}{1 + \varepsilon} (1 - \varepsilon \tau t)\right) \frac{\sin(\tau t)}{\tau}
\le \varepsilon (2 + \tau t) \frac{\sin(\tau t)}{\tau}
< 3\varepsilon t,$$

from which (51) follows, as $\sin(\tau t) \leq \frac{26}{29}\tau t$ for $t > \frac{1}{\tau}$. Then, combining (50) and (51), we obtain (49), which proves the lemma.

We next wish to prove the following estimate.

Lemma 2.4. If $\beta(t)$ satisfies (3) and (5), $k \leq 1$, $\frac{\pi - \varepsilon}{k} \leq \nu \leq \frac{\pi}{k}$ and $\varepsilon \leq \frac{1}{5}$, then

(52)
$$\phi(\sigma, \tau) \ge \frac{1}{2}\widehat{\beta}(\sigma).$$

Proof. Beginning with the formulas

$$\phi(\sigma,\tau) = \int_0^\infty \cos(\tau t) e^{-\sigma t} c(t) dt + \frac{\sigma \beta(\infty)}{\sigma^2 + \tau^2} \quad \text{and} \quad \widehat{\beta}(\sigma) = \int_0^\infty e^{-\sigma t} c(t) dt + \frac{\beta(\infty)}{\sigma},$$

we note that as $\frac{\pi-\varepsilon}{k} \leq \nu \leq \frac{\pi}{k}$ gives us $\frac{2-\varepsilon}{k} \leq \sigma \leq \frac{2}{k}$ and $0 \leq \tau \leq \frac{\varepsilon}{k}$, we find that $\frac{\sigma}{\sigma^2+\tau^2} \geq \frac{(2-\varepsilon)k}{4+\varepsilon^2} \geq \frac{45k}{101} > \frac{k}{4-2\varepsilon} \geq \frac{1}{2\sigma}$, so we need only show that $\phi_c(\sigma,\tau) \geq \frac{1}{2}\widehat{c}(\sigma)$. Integrating by parts, we obtain

$$\phi_c(\sigma, \tau) = \frac{1}{\sigma^2 + \tau^2} \int_0^\infty [\sigma(1 - e^{-\sigma t} \cos(\tau t)) + \tau e^{-\sigma t} \sin(\tau t)] (-c'(t)) dt$$

and

$$\widehat{c}(\sigma) = \frac{1}{\sigma} \int_0^\infty (1 - e^{-\sigma t})(-c'(t)) dt.$$

The boundary terms vanish as in the proof of Lemma 2.3. Then, setting

$$f(t) = \frac{1}{\sigma^2 + \tau^2} [\sigma(1 - e^{-\sigma t}\cos(\tau t)) + \tau e^{-\sigma t}\sin(\tau t)]$$
 and $g(t) = \frac{1}{\sigma}(1 - e^{-\sigma t}),$

we see that for $t \leq \frac{1}{\tau}$, we have $f'(t) = e^{-\sigma t}\cos(\tau t) \geq \frac{1}{2}e^{-\sigma t} = \frac{1}{2}g'(t)$, so as f(0) = g(0) = 0, we may integrate over [0,t], for $t \leq \frac{1}{\tau}$, to obtain $f(t) \geq \frac{1}{2}g(t)$ for $t \leq \frac{1}{\tau}$. Thus, it is only a matter of showing that $f(t) \geq \frac{1}{2}g(t)$ for $t > \frac{1}{\tau}$ to establish (52). As $\frac{2-\varepsilon}{k} \leq \sigma \leq \frac{2}{k}$ and $0 \leq \tau \leq \frac{\varepsilon}{k}$, it follows that for $\varepsilon \leq \frac{1}{5}$ we have $\frac{\sigma}{\tau} \geq \frac{2-\varepsilon}{\varepsilon} \geq 9$; so clearly we have $e^{\frac{-\sigma}{\tau}} \leq (1-e^{-\sigma t})$ for $t > \frac{1}{\tau}$. Thus, for $t > \frac{1}{\tau}$, we see that

$$\begin{split} f(t) &\geq \frac{k^2}{4+\varepsilon^2} \left(\frac{2-\varepsilon}{k} (1-e^{-\sigma t}) - \frac{\varepsilon}{k} e^{\frac{-\sigma}{\tau}} \right) \\ &\geq \frac{k}{4+\varepsilon} \left((2-\varepsilon)(1-e^{-\sigma t}) - \varepsilon(1-e^{-\sigma t}) \right) \\ &= k \left(\frac{2-2\varepsilon}{4+\varepsilon} \right) (1-e^{-\sigma t}) \\ &\geq \frac{8k}{21} (1-e^{-\sigma t}) \\ &\geq \frac{1}{2} g(t), \end{split}$$

which proves the lemma. \Box

We next wish to extend Lemma 3.1 in [21].

LEMMA 2.5. If $\beta(t)$ satisfies (3) and (5), $\lambda \geq \lambda_0$, and k < 1, then

(53)
$$k \sum_{n=1}^{\infty} |U^n(\lambda)| \le C\lambda.$$

Proof. Following [21], we define the generating function of $\{U^n(\lambda)\}_{n=0}^{\infty}$ to be

$$\widetilde{U}(z,\lambda) = \sum_{n=1}^{\infty} U^n(\lambda) z^n,$$

which may be shown to obey the relations

$$\widetilde{U}(z,\lambda) = \frac{z}{k} \frac{1}{\left(\frac{1-z}{k}\right) + \lambda \widehat{\beta}\left(\frac{1-z}{k}\right)} = \frac{z}{k} \widehat{u}\left(\frac{1-z}{k},\lambda\right).$$

Then, an application of Hardy's inequality [4, p. 48] gives us

$$\sum_{n=1}^{\infty} |U^{n}(\lambda)| \leq 2 \sum_{n=1}^{\infty} \frac{n|U^{n}(\lambda)|}{n+1} \leq \int_{-\pi}^{\pi} |\widetilde{U}'(e^{i\nu}, \lambda)| \, d\nu$$

$$\leq 2k \int_{0}^{\frac{\pi}{k}} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu = 2k \left\{ \int_{0}^{\varepsilon} + \int_{\varepsilon}^{\frac{\varepsilon}{k}} + \int_{\frac{\varepsilon}{k}}^{\frac{\pi}{k}} \right\} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu,$$

where

$$\widetilde{U}'(z,\lambda) = \frac{1}{k} \frac{1}{\left\lceil \left(\frac{1-z}{k}\right) + \lambda \widehat{\beta} \left(\frac{1-z}{k}\right) \right\rceil^2} \left[\frac{1}{k} + \lambda \widehat{\beta} \left(\frac{1-z}{k}\right) + \frac{\lambda z}{k} \widehat{\beta}' \left(\frac{1-z}{k}\right) \right].$$

Thus, our extension reduces to establishing the three estimates

(55)
$$k^2 \int_0^{\varepsilon} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu \le C\lambda^{-1},$$

(56)
$$k^2 \int_{\varepsilon}^{\frac{\varepsilon}{k}} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu \le C\lambda,$$

(57)
$$k^2 \int_{\frac{\varepsilon}{L}}^{\frac{\pi}{k}} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu \le Ck\lambda.$$

We prove (55) first. For $\varepsilon, k < 1$, we see that when $0 \le \nu \le \varepsilon$, we get

(58)
$$\frac{\nu}{2} \le \tau(k,\nu) \le \nu, \quad \sigma(k,\nu) \le \varepsilon \tau \le \tau, \quad \cos(k\nu) \ge \frac{1}{2}, \quad \frac{\sin(k\nu)}{k} \le \varepsilon.$$

By (40) and (58), we see that

(59)
$$\operatorname{Re}\widehat{\beta}(\sigma + i\tau) \ge C \int_0^{\frac{1}{\tau}} -te^{-\sigma t}\beta'(t) dt \ge C \int_0^{\frac{1}{\varepsilon}} -t\beta'(t) dt \ge \sqrt{2}C \frac{\tau}{\lambda}.$$

We may obviously assume C < 1 in (59), giving

$$\begin{split} |D(\sigma+i\tau,\lambda)|^2 &= \left|\phi(\sigma,\tau) + \frac{\sigma}{\lambda}\right|^2 + \left|\tau\theta(\sigma,\tau) - \frac{\tau}{\lambda}\right|^2 \\ &\geq \phi^2(\sigma,\tau) + C^2 \left|\tau\theta(\sigma,\tau) - \frac{\tau}{\lambda}\right|^2 \\ &\geq \frac{\phi^2(\sigma,\tau)}{2} + C^2 \frac{\tau^2}{\lambda^2} + C^2 \left|\tau\theta(\sigma,\tau) - \frac{\tau}{\lambda}\right|^2 \\ &= \frac{\phi^2(\sigma,\tau)}{2} + C^2 \frac{\tau^2}{2} \left(\theta(\sigma,\tau) - \frac{2}{\lambda}\right)^2 + C^2 \frac{\tau^2}{2} \theta^2(\sigma,\tau) \\ &\geq C^2 \left(\frac{\phi^2(\sigma,\tau) + \tau^2\theta^2(\sigma,\tau)}{2} + \frac{\tau^2}{2} \left(\theta(\sigma,\tau) - \frac{2}{\lambda}\right)^2\right) \\ &\geq C^2 \left(\frac{\phi^2(\sigma,\tau) + \tau^2\theta^2(\sigma,\tau)}{2}\right), \end{split}$$

so we obtain

(60)
$$|D(\sigma + i\tau, \lambda)| \ge C|\widehat{\beta}(\sigma + i\tau)|.$$

Then, (36), (42), and (58) give us

$$|\widehat{\beta}(\sigma+i\tau)| \ge C \int_0^{\frac{1}{\tau}} e^{-\sigma t} \beta(t) \, dt \ge C|\widetilde{\beta}(\tau)| \ge C \int_0^{\frac{1}{\varepsilon}} \beta(t) \, dt \ge C|\widetilde{\beta}(\varepsilon)|,$$

so (60) implies that

(61)
$$|D(\sigma + i\tau, \lambda)| \ge C|\widetilde{\beta}(\tau)| \ge C|\widetilde{\beta}(\varepsilon)|.$$

Note that (36) and (42) give us

(62)
$$|\widehat{\beta}(\sigma + i\tau)| = |(e^{-\sigma t}\beta(t))^{\sim}(\tau)| \le C|\widetilde{\beta}(\tau)|, \qquad \tau > 0$$

We also see that (37) and (43) imply

(63)
$$|\widehat{\beta}'(\sigma + i\tau)| = \left| \frac{d}{d\tau} (e^{-\sigma t} \beta(t)) \widetilde{}(\tau) \right| \le C\theta(\tau), \qquad \sigma, \tau > 0.$$

Then it follows from (36), (37), (44), (61), (62), and (63) that

$$\begin{split} k^2 \int_0^\varepsilon |\widetilde{U}'(e^{-ik\nu},\lambda)| \, d\nu &\leq \int_0^\varepsilon \frac{1}{\lambda^2 |D^2(\sigma+i\tau,\lambda)|} \left[1 + \lambda k |\widehat{\beta}(\sigma+i\tau)| + \lambda |\widehat{\beta}'(\sigma+i\tau)| \right] \, d\nu \\ &\leq \frac{C}{\lambda} \int_0^\varepsilon \frac{1}{|\widetilde{\beta}(\varepsilon)|^2} + \frac{k |\widetilde{\beta}(\tau)|}{|\widetilde{\beta}(\tau)| |\widetilde{\beta}(\varepsilon)|} + \frac{\theta(\tau)}{|\widetilde{\beta}(\tau)|^2} \, d\tau \\ &\leq C\lambda^{-1}, \end{split}$$

so estimate (55) holds. We now show (56). Here $\varepsilon \leq \nu \leq \frac{\varepsilon}{k}$ and for $\varepsilon, k \leq 1$ we have

(64)
$$\frac{\nu}{2} \le \tau(k,\nu) \le \nu, \quad \sigma(k,\nu) \le \varepsilon \tau(k,\nu) \le \tau(k,\nu), \quad \cos(k\nu) \ge \frac{1}{2}.$$

We shall establish the following estimates on $|D(\sigma+i\tau,\lambda)|$ when $\varepsilon \leq \min\left\{\frac{C_1}{1392}, \frac{5}{348}\right\}$:

$$(65) |D(\sigma + i\tau, \lambda)| \ge C \left(\tau \theta(\tau) + \int_0^{\frac{1}{\tau}} \beta(t) \, dt \right), \tau \in \left[\frac{\varepsilon}{2}, \frac{\omega}{2} \right] \bigcap \left[\frac{\varepsilon}{2}, \frac{\varepsilon}{k} \right],$$

and

(66)
$$|D(\sigma + i\tau, \lambda)| \ge C \frac{\tau - \omega}{\lambda}, \qquad \tau \ge 2\omega.$$

We show (65) first. To establish

(67)
$$|D(\sigma+i\tau,\lambda)| \ge C\tau\theta(\tau), \qquad \tau \in \left[\frac{\varepsilon}{2},\frac{\omega}{2}\right] \bigcap \left[\frac{\varepsilon}{2},\frac{\varepsilon}{k}\right],$$

note that in the case where

$$\tau \left| \theta(\tau) - \frac{1}{\lambda} \right| < \frac{C_1 \tau}{2} \int_0^{\frac{1}{\tau}} t \beta(t) dt$$

we can use (37), (41), (45), and (64) to show that

$$|D(\sigma + i\tau, \lambda)| \ge C\phi(\tau) \ge C\left(C_1\tau \int_0^{\frac{1}{\tau}} t\beta(t) dt - \tau \left|\theta(\tau) - \frac{1}{\lambda}\right|\right)$$

$$\ge C\left(\frac{C_1\tau}{2} \int_0^{\frac{1}{\tau}} t\beta(t) dt\right) \ge C\tau\theta(\tau).$$

Similarly, when

$$\tau \left| \theta(\tau) - \frac{1}{\lambda} \right| \ge \frac{C_1 \tau}{2} \int_0^{\frac{1}{\tau}} t \beta(t) dt,$$

we find that (37) and (48) give us

$$\tau \left| \theta(\tau) - \frac{1}{\lambda} \right| \ge \frac{C_1 \tau}{2} \int_0^{\frac{1}{\tau}} t\beta(t) dt \ge \frac{C_1 \tau}{24} \theta(\tau) \ge 2\tau |\theta(\sigma, \tau) - \theta(\tau)|.$$

Thus, it follows by (37) that

$$\begin{split} |D(\sigma+i\tau,\lambda)| &\geq \tau \left| \theta(\sigma,\tau) - \frac{1}{\lambda} \right| = \tau \left| \left(\theta(\tau) - \frac{1}{\lambda} \right) + \left(\theta(\sigma,\tau) - \theta(\tau) \right) \right| \geq \frac{\tau}{2} \left| \theta(\tau) - \frac{1}{\lambda} \right| \\ &\geq \frac{C_1 \tau}{2} \int_0^{\frac{1}{\tau}} t \beta(t) \, dt \geq C \tau \theta(\tau). \end{split}$$

This establishes (67). The estimate

(68)
$$|D(\sigma + i\tau, \lambda)| \ge C \int_0^{\frac{1}{\tau}} \beta(t) dt, \qquad \tau \in \left[\frac{\varepsilon}{2}, \frac{\omega}{2}\right] \bigcap \left[\frac{\varepsilon}{2}, \frac{\varepsilon}{k}\right],$$

follows from the same argument with $\tau \int_0^{\frac{1}{\tau}} t\beta(t) dt$ replaced by $\int_0^{\frac{1}{\tau}} \beta(t) dt$. Then, combining (67) and (68), we have (65). To prove (66), we note that (41) gives us

$$|D(\sigma + i\tau, \lambda)| \ge C\phi(\tau),$$

which establishes the estimate when $\phi(\tau) \geq \frac{C_1}{2} \frac{\tau - \omega}{\lambda}$. Thus, assume $\phi(\tau) < \frac{C_1}{2} \frac{\tau - \omega}{\lambda}$. Then, (45) gives us

$$C_1 \frac{\tau - \omega}{\lambda} \le |D(i\tau, \lambda)| \le \phi(\tau) + \tau \left| \theta(\tau) - \frac{1}{\lambda} \right| \le \frac{C_1}{2} \frac{\tau - \omega}{\lambda} + \tau \left| \theta(\tau) - \frac{1}{\lambda} \right|,$$

SO

(69)
$$\left| \frac{1}{\lambda} - \theta(\tau) \right| \ge \frac{C_1(\tau - \omega)}{2\tau\lambda}, \qquad \tau \ge 2\omega.$$

As $\theta(\tau) = \theta_c(\tau) + \frac{\beta(\infty)}{\tau^2}$, we see by (4.4) of [1] that $\theta(\tau)$ is decreasing. It follows by our construction of $\omega(\lambda)$ that in the case where $\lambda \geq \lambda_1$, or by (69), in the case where $\lambda_0 \geq \lambda > \lambda_1$ with $\frac{1}{\lambda} - \theta(\tau) \geq \frac{C_1(\tau - \omega)}{2\tau\lambda}$, that as $\varepsilon \leq \min\left\{\frac{C_1}{1392}, \frac{5}{348}\right\}$, we have

$$\begin{aligned} \theta(\sigma,\tau) &\leq (1+29\varepsilon)\theta(\tau) = (1+29\varepsilon)\frac{1}{\lambda} - (1+29\varepsilon)\left(\frac{1}{\lambda} - \theta(\tau)\right) \\ &\leq \frac{1}{\lambda} - \frac{C_1(\tau-\omega)}{4\tau\lambda}\left(2+58\varepsilon - \frac{116\varepsilon\tau}{C_1(\tau-\omega)}\right) \\ &\leq \frac{1}{\lambda} - \frac{C_1(\tau-\omega)}{4\tau\lambda}\left(2 - \frac{232\varepsilon}{C_1}\right) \\ &\leq \frac{1}{\lambda} - \frac{C_1(\tau-\omega)}{4\tau\lambda}, \end{aligned}$$

so we find that

(70)
$$|D(\sigma + i\tau, \lambda)| \ge \tau \left(\frac{1}{\lambda} - \theta(\sigma, \tau)\right) \ge C \frac{\tau - \omega}{\lambda}.$$

Similarly, by (69), in the case where $\lambda_0 \geq \lambda > \lambda_1$ with $\frac{1}{\lambda} - \theta(\tau) < -\frac{C_1(\tau - \omega)}{2\tau\lambda}$, we see that as $\varepsilon \leq \min\left\{\frac{C_1}{1392}, \frac{5}{348}\right\}$ and as $\frac{\tau}{\tau - \omega} \leq 2$ for $\tau \geq 2\omega$, we have

$$\theta(\sigma,\tau) \ge (1 - 29\varepsilon)\theta(\tau) = (1 - 29\varepsilon)\frac{1}{\lambda} + (1 - 29\varepsilon)\left(\theta(\tau) - \frac{1}{\lambda}\right)$$

$$\ge \frac{1}{\lambda} + \frac{C_1(\tau - \omega)}{4\tau\lambda}\left(2 - 58\varepsilon - \frac{116\varepsilon\tau}{C_1(\tau - \omega)}\right)$$

$$\ge \frac{1}{\lambda} + \frac{C_1(\tau - \omega)}{4\tau\lambda}\left(2 - 58\varepsilon - \frac{232\varepsilon}{C_1}\right)$$

$$\ge \frac{1}{\lambda} + \frac{C_1(\tau - \omega)}{4\tau\lambda},$$

so we obtain

(71)
$$|D(\sigma + i\tau, \lambda)| \ge \tau \left(\theta(\sigma, \tau) - \frac{1}{\lambda}\right) \ge C \frac{\tau - \omega}{\lambda}.$$

Combining (70) and (71) completes the proof of (66).

Next, let $E_1 = \left[\frac{\varepsilon}{2}, \frac{\omega}{2}\right] \cap \left[\frac{\varepsilon}{2}, \frac{\varepsilon}{k}\right], E_2 = \left[\frac{\omega}{2}, 2\omega\right] \cap \left[\frac{\varepsilon}{2}, \frac{\varepsilon}{k}\right],$ and $E_3 = \left[2\omega, \infty\right) \cap \left[\frac{\varepsilon}{2}, \frac{\varepsilon}{k}\right].$ Note that our construction of ω and the decreasing nature of $\theta(\tau)$ give us $\lambda\theta(\tau) \geq C$ for $\tau \in E_1$. Also, by (36), we see that $|\widetilde{\beta}(\tau)| \leq C$ for $\tau \in E_3$. Then by (17), (36), (41), (48), (62), (63), (65), and (66), we see that

$$\begin{split} k^2 \int_{\varepsilon}^{\frac{\varepsilon}{k}} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu \\ & \leq C \left\{ \int_{E_1} + \int_{E_2} + \int_{E_3} \right\} \frac{1}{\lambda^2 |D^2(\sigma + i\tau, \lambda)|} \left[1 + \lambda k |\widehat{\beta}(\sigma + i\tau)| + \lambda |\widehat{\beta}'(\sigma + i\tau)| \right] \, d\tau \\ & \leq C \left(\int_{\frac{\varepsilon}{2}}^{\frac{\omega}{2}} \frac{d\tau}{\tau^2} + k \int_{E_1} \frac{A(\tau^{-1})}{\lambda \tau \theta(\tau) A(\tau^{-1})} d\tau \right) \\ & \quad + \frac{C}{\lambda} \left(\int_{E_2} \frac{d\tau}{\lambda \phi^2(\tau)} + k \frac{\phi(\tau) + \tau \theta(\tau)}{\phi^2(\tau)} + \frac{\theta(\tau)}{\phi^2(\tau)} d\tau \right) \\ & \quad + C \left(\int_{E_3} \frac{1 + \lambda (1 + \theta(\tau))}{(\tau - \omega)^2} \, d\tau \right), \end{split}$$

so by (46) and (47) we find that

$$k^2 \int_{\varepsilon}^{\frac{\varepsilon}{k}} |\widetilde{U}'(e^{-ik\nu},\lambda)| \, d\nu \leq C(1+\lambda+k\lambda) \leq C\lambda,$$

which establishes (56). We now establish (57). Here $\frac{\varepsilon}{k} \leq \nu \leq \frac{\pi}{k}$, and we have

(72)
$$\sigma(k,\nu) \ge \frac{1 - \cos(\varepsilon)}{k} \equiv \frac{C(\varepsilon)}{k}, \qquad |D(\sigma + i\tau)| \ge \frac{\sigma}{\lambda} \ge \frac{C(\varepsilon)}{k\lambda}.$$

Then, as $k \leq 1$, we have

$$(73) \qquad |\widehat{\beta}(\sigma+i\tau)| = \left| \int_0^\infty e^{-(\sigma+i\tau)t} \beta(t) \, dt \right| \le \int_0^\infty e^{-C(\varepsilon)t} \beta(t) \, dt = \widehat{\beta}(C(\varepsilon))$$

and

$$(74) \qquad |\widehat{\beta}'(\sigma+i\tau)| = \left| \int_0^\infty -te^{-(\sigma+i\tau)t} \beta(t) \, dt \right| \le \int_0^\infty te^{-C(\varepsilon)t} \beta(t) \, dt = |\widehat{\beta}'(C(\varepsilon))|.$$

So, (72), (73), and (74) give us

$$\begin{split} k^2 \int_{\frac{\pi}{k}}^{\frac{\pi}{k}} |\widetilde{U}'(e^{-ik\nu}, \lambda)| \, d\nu &\leq \int_{\frac{\pi}{k}}^{\frac{\pi}{k}} \frac{1}{\lambda^2 |D^2(\sigma + i\tau, \lambda)|} \left[1 + \lambda k |\widehat{\beta}(\sigma + i\tau)| + \lambda |\widehat{\beta}'(\sigma + i\tau)| \right] \, d\nu \\ &\leq \int_{\frac{\pi}{k}}^{\frac{\pi}{k}} \frac{k^2}{(C(\varepsilon))^2} \left(1 + \lambda k \widehat{\beta}(C(\varepsilon)) + \lambda |\widehat{\beta}'(C(\varepsilon))| \right) \\ &\leq Ck\lambda, \end{split}$$

so estimate (57) holds. This proves the lemma.

3. Proof of Theorems 1 and 2. Here we adopt the overall strategy of Xu in proving our theorems, and we refer the reader to [21] for the preliminaries of the proof. We remark that Xu establishes the formula

(75)
$$U^{n}(\lambda) = \operatorname{Re}\left\{\frac{1}{\pi t_{n-1}\lambda} \int_{0}^{\frac{\pi}{k}} e^{i\nu t_{n-2}} \frac{D_{s}(s(k,\nu),\lambda)}{D^{2}(s(k,\nu),\lambda)} d\nu\right\}.$$

Then, following [1], this integral is decomposed into the five parts:

(76)
$$U^{n}(\lambda) = \operatorname{Re}\{\lambda^{-1}U_{1}^{n} + \lambda^{-2}U_{2}^{n} + \lambda^{-3}U_{3}^{n} + U_{4}^{n}(\lambda) + U_{5}^{n}(\lambda)\},\$$

where

(77)
$$U_1^n = \frac{1}{\pi t_{n-1}} \int_0^{\varepsilon} e^{i\nu t_{n-2}} \frac{\widehat{\beta}'(s)}{[\widehat{\beta}(s)]^2} d\nu,$$

(78)
$$U_2^n = \frac{1}{\pi t_{n-1}} \int_0^{\varepsilon} e^{i\nu t_{n-2}} \frac{1}{[\widehat{\beta}(s)]^2} \left(1 - \frac{2s\widehat{\beta}'(s)}{\widehat{\beta}(s)} \right) d\nu,$$

(79)
$$U_3^n = \frac{-1}{\pi t_{n-1}} \int_0^\varepsilon e^{i\nu t_{n-2}} \frac{2s}{[\widehat{\beta}(s)]^3} d\nu,$$

(80)
$$U_4^n(\lambda) = \frac{1}{\pi t_{n-1\lambda^3}} \int_0^\varepsilon e^{i\nu t_{n-2}} \frac{s^2 D_s(s,\lambda)}{[\widehat{\beta}(s)]^2 D(s,\lambda)} \left(\frac{2}{\widehat{\beta}(s)} + \frac{1}{D(s,\lambda)}\right) d\nu,$$

(81)
$$U_5^n(\lambda) = \frac{1}{\pi t_{n-1\lambda}} \int_{\varepsilon}^{\frac{\pi}{k}} e^{i\nu t_{n-2}} \frac{D_s(s,\lambda)}{[D(s,\lambda)]^2} d\nu.$$

For m > 2, we will establish the estimates

(82)
$$|U_4^n(\lambda)| \le Ct_{n-2}^{-2}, \qquad n \ge m, \qquad \lambda \ge \lambda_0$$

and either

(83)
$$|U_5^n(\lambda)| \le Ct_{n-2}^{-2}\lambda, \qquad n \ge m, \qquad \lambda \ge \lambda_0$$

or

(84)
$$|U_5^n(\lambda)| \le Ct_{n-2}^{-2}, \qquad n \ge m, \qquad \lambda \ge \lambda_0$$

to prove Theorem 1 or 2, respectively. This follows, as then we could insert three different values of λ into (76) and, by utilizing (53), solve for the U_j^n (j = 1, 2, 3) to show

$$k \sum_{n=m+1}^{\infty} (|\text{Re}\{U_1^n\}| + |\text{Re}\{U_2^n\}| + |\text{Re}\{U_3^n\}|) < \infty.$$

Then, (34) or (35) follow, giving us Theorem 1 and 2, respectively. We show (82) first. Integrating (80) and (81) by parts, we obtain

$$U_4^n(\lambda) = \frac{1}{i\pi t_{n-2}t_{n-1}\lambda^3} e^{i\nu t_{n-2}} \frac{s^2 D_s(s,\lambda)}{[\widehat{\beta}(s)]^2 D(s,\lambda)} \left(\frac{2}{\widehat{\beta}(s)} + \frac{1}{D(s,\lambda)}\right) \Big|_{\nu=0}^{\varepsilon}$$

$$- \frac{1}{\pi t_{n-2}t_{n-1}\lambda^3} \int_0^{\varepsilon} e^{i\nu t_{n-3}} \left[\frac{2s D_s(s,\lambda) + s^2 \widehat{\beta}''(s)}{[\widehat{\beta}(s)]^2 D(s,\lambda)} \left(\frac{2}{\widehat{\beta}(s)} + \frac{1}{D(s,\lambda)}\right) - \frac{s^2 D_s(s,\lambda)}{[\widehat{\beta}(s)]^2 D(s,\lambda)} \right]$$

$$- \frac{s^2 D_s(s,\lambda)}{[\widehat{\beta}(s)]^2} + \frac{4\widehat{\beta}'(s) + 2\lambda^{-1}}{\widehat{\beta}(s) D(s,\lambda)} + \frac{2D_s(s,\lambda)}{D^2(s,\lambda)} \right] d\nu$$

$$(85)$$

and

$$(86) U_5^n(\lambda) = \frac{1}{i\pi t_{n-2}t_{n-1}\lambda} e^{i\nu t_{n-2}} \frac{D_s(s,\lambda)}{D^2(s,\lambda)} \Big|_{\nu=\varepsilon}^{\frac{\pi}{k}} - \frac{1}{i\pi t_{n-2}t_{n-1}\lambda} \int_{\varepsilon}^{\frac{\pi}{k}} e^{i\nu t_{n-3}} \left[\frac{\widehat{\beta}''(s)}{D^2(s,\lambda)} - \frac{2D_s^2(s,\lambda)}{D^3(s,\lambda)} \right] d\nu.$$

We see by (41) and (48) that for $0 < \nu \le \varepsilon$ with ε appropriately small

$$(87) \qquad |\widehat{\beta}(s)| = \sqrt{\phi^2(\sigma,\tau) + \tau^2\theta^2(\sigma,\tau)} \geq C\sqrt{\phi^2(\tau) + \tau^2\theta^2(\tau)} = C|\widehat{\beta}(i\tau)|.$$

Then, (36), (37), (58), (63), and (87) imply that the boundary term in (85) vanishes at $\nu = 0$. Note that [1, eq. (5.3)] and (37) give us, for $\sigma, \tau > 0$,

$$|\widehat{\beta}''(s)| = \left| \frac{d^2}{d\tau^2} (e^{-\sigma t} \beta(t)) \widetilde{}(\tau) \right| \le C \int_0^{\frac{1}{\tau}} t^2 e^{-\sigma t} \beta(t) dt$$

$$\le \frac{C}{\tau} \int_0^{\frac{1}{\tau}} t e^{-\sigma t} \beta(t) dt \le \frac{C}{\tau} \theta(\tau).$$
(88)

Then, (36), (37), (58), (63), (87), and (88) allow us to establish the estimate (82) in a manner similar to (5.6) of [21].

To establish (83) and (84), we first note that as $\int_0^{\frac{2}{\sigma}} te^{-\sigma t} dt \ge \int_{\frac{2}{\sigma}}^{\infty} te^{-\sigma t} dt$ for all $\sigma > 0$, it follows from the decreasing nature of $\beta(t)$ and (36) that

$$|\widehat{\beta}'(\sigma)| = \left(\int_0^{\frac{2}{\sigma}} + \int_{\frac{2}{\sigma}}^{\infty}\right) t e^{-\sigma t} \beta(t) dt \le 2 \int_0^{\frac{2}{\sigma}} t e^{-\sigma t} \beta(t) dt$$

$$\le \frac{4}{\sigma} \int_0^{\frac{2}{\sigma}} e^{-\sigma t} \beta(t) dt \le \frac{C}{\sigma} \widehat{\beta}(\sigma).$$
(89)

Also, as $\int_0^{\frac{1}{\tau}} e^{-\tau t} dt \ge \int_{\frac{1}{\tau}}^{\infty} e^{-\tau t} dt$ for all $\tau > 0$, we see by the decreasing nature of $\beta(t)$, (36), (42), (58), and (60) that, for $0 < \sigma \le \varepsilon \tau < \tau$,

$$|D(s,\lambda)| \ge C|\widehat{\beta}(\sigma+i\tau)| \ge C \int_0^{\frac{1}{\tau}} e^{-\sigma t} \beta(t) \, dt \ge C \int_0^{\frac{1}{\tau}} e^{-\tau t} \beta(t) \, dt$$

$$\ge C_0 \left(\int_0^{\frac{1}{\tau}} + \int_{\frac{1}{\tau}}^{\infty} \right) e^{-\tau t} \beta(t) \, dt = C\widehat{\beta}(\tau).$$

Then, by (89) and (90), we are able to estimate the boundary terms in (86) as on pp. 148–149 of [21]. From this we obtain

$$(91) \quad |U_5^n(\lambda)| \le Ct_{n-2}^{-2} \left[k^2 + \lambda^{-1} + \lambda^{-1} \left(\int_{\varepsilon}^{\frac{\pi}{k}} \frac{|\widehat{\beta}''(s)|}{|D^2(s,\lambda)|} + \frac{|D_s^2(s,\lambda)|}{|D^3(s,\lambda)|} d\nu \right) \right].$$

We decompose the interval of integration into the three intervals $E_1 = \left[\varepsilon, \frac{\varepsilon}{k}\right), E_2 = \left[\frac{\varepsilon}{k}, \frac{\pi-\varepsilon}{k}\right)$, and $E_3 = \left[\frac{\pi-\varepsilon}{k}, \frac{\pi}{k}\right)$. To estimate the integral on E_1 , we note that when $\nu \in E_1$, we have $0 < \sigma \le \varepsilon \tau < \tau$. So by (17), (37), (41), (47), (63), (65), (66), and (88), it follows that

$$\lambda^{-1} \int_{\varepsilon}^{\frac{\varepsilon}{k}} \left(\frac{|\widehat{\beta}''(s)|}{|D(s,\lambda)|^{2}} + \frac{|D_{s}(s,\lambda)|^{2}}{|D(s,\lambda)|^{3}} \right) d\nu$$

$$\leq C\lambda^{-1} \left(\int_{\frac{\varepsilon}{2}}^{\frac{\omega}{2}} \frac{\lambda^{-2} + \theta^{2}(\tau)}{\tau^{3}\theta^{3}(\tau)} d\tau + \int_{\frac{\omega}{2}}^{2\omega} \frac{\lambda^{-1}}{\tau\phi^{2}(\tau)} d\tau + \int_{\frac{\omega}{2}}^{2\omega} \frac{\lambda^{-2}}{\phi^{3}(\tau)} d\tau \right)$$

$$+ \int_{2\omega}^{\infty} \frac{\lambda d\tau}{\tau |\tau - \omega|^{2}} + \int_{2\omega}^{\infty} \frac{\lambda d\tau}{|\tau - \omega|^{3}} d\tau$$
(92)

Then, by (37), (46), (47), and either (17) or (18), we obtain

(93)
$$\lambda^{-1} \int_{\varepsilon}^{\frac{\varepsilon}{k}} \left(\frac{|\widehat{\beta}''(s)|}{|D(s,\lambda)|^2} + \frac{|D_s(s,\lambda)|^2}{|D(s,\lambda)|^3} \right) d\nu \le C(\lambda+1)$$

or

(94)
$$\lambda^{-1} \int_{\varepsilon}^{\frac{\varepsilon}{k}} \left(\frac{|\widehat{\beta}''(s)|}{|D(s,\lambda)|^2} + \frac{|D_s(s,\lambda)|^2}{|D(s,\lambda)|^3} \right) d\nu \le C,$$

respectively.

To estimate the integrals on E_2 and E_3 , we first note that as $\int_0^{\frac{3}{\sigma}} t^2 e^{-\sigma t} dt \ge \int_{\frac{3}{\sigma}}^{\infty} t^2 e^{-\sigma t} dt$ for all $\sigma > 0$, the decreasing nature of $\beta(t)$ gives us, for $\sigma > 0$,

$$|\widehat{\beta}''(s)| = \left| \int_0^\infty t^2 e^{-i\tau t} (e^{-\sigma t} \beta(t)) dt \right| \le \int_0^\infty t^2 e^{-\sigma t} \beta(t) dt = \widehat{\beta}''(\sigma)$$

$$= \left(\int_0^{\frac{3}{\sigma}} + \int_{\frac{3}{\sigma}}^\infty \right) t^2 e^{-\sigma t} \beta(t) dt \le 2 \int_0^{\frac{3}{\sigma}} t^2 e^{-\sigma t} \beta(t) dt$$

$$\le \frac{18}{\sigma^2} \int_0^{\frac{3}{\sigma}} e^{-\sigma t} \beta(t) dt \le \frac{C}{\sigma^2} \widehat{\beta}(\sigma).$$
(95)

Then, in the case where $\nu \in E_2$, we see that $\sigma \leq \frac{2}{k}$ and $\tau \geq \frac{\sin(\varepsilon)}{k}$. Then (39) and (40) give, for $\nu \in E_2$,

$$\phi(\sigma,\tau) \ge -C \int_0^{\frac{1}{\tau}} t e^{-\sigma t} \beta'(t) dt \ge -C e^{\frac{-\sigma}{\tau}} \int_0^{\frac{1}{\tau}} t \beta'(t) dt$$

$$\ge -C e^{\frac{-2}{\sin(\varepsilon)}} \int_0^{\frac{1}{\tau}} t \beta'(t) dt \ge C \phi(\tau).$$

So, we see that (17), (63), (88), and (96) give us

$$\lambda^{-1} \int_{\frac{\varepsilon}{k}}^{\frac{\pi-\varepsilon}{k}} \frac{|\widehat{\beta}''(s)|}{|D^{2}(s,\lambda)|} + \frac{|D_{s}^{2}(s,\lambda)|}{|D^{3}(s,\lambda)|} d\nu$$

$$\leq C\lambda^{-1} \int_{\frac{\varepsilon}{k}}^{\frac{\pi-\varepsilon}{k}} \frac{k\lambda\theta(\tau)}{\tau\phi(\tau)} + k^{3}\lambda + \frac{k^{2}\lambda\theta(\tau)}{\phi(\tau)} + \frac{k\lambda\theta^{2}(\tau)}{\phi^{2}(\tau)} d\nu$$

$$\leq C.$$
(97)

For the integral on E_3 , note that since we have $\frac{2-\varepsilon}{k} \le \sigma \le \frac{2}{k}$ and $0 \le \tau \le \frac{\varepsilon}{k}$ for $\nu \in E_3$, we see by (52), (89), and (95) that

$$\lambda^{-1} \int_{\frac{\pi-\varepsilon}{k}}^{\frac{\pi}{k}} \frac{|\widehat{\beta}''(s)|}{|D^{2}(s,\lambda)|} + \frac{|D_{s}^{2}(s,\lambda)|}{|D^{3}(s,\lambda)|} d\nu$$

$$\leq C\lambda^{-1} \int_{\frac{\pi-\varepsilon}{k}}^{\frac{\pi}{k}} \frac{\lambda \widehat{\beta}(\sigma)}{\sigma^{3} \widehat{\beta}(\sigma)} + \frac{\lambda}{\sigma^{3}} + \frac{\lambda \widehat{\beta}(\sigma)}{\sigma^{3} \widehat{\beta}(\sigma)} + \frac{\lambda [\widehat{\beta}(\sigma)]^{2}}{\sigma^{3} [\widehat{\beta}(\sigma)]^{2}} d\nu$$
(98)
$$\leq C.$$

Then, by (97), (98), and either (93) or (94), we have established (83) or (84), respectively, and thus we have proven Theorems 1 or 2, respectively.

REFERENCES

- R. W. CARR AND K. B. HANNSGEN, A nonhomogeneous integrodifferential equation in Hilbert space, SIAM J. Math. Anal., 10 (1979), pp. 961–984.
- [2] R. W. CARR AND K. B. HANNSGEN, Resolvent formulas for a Volterra equation in Hilbert space, SIAM J. Math. Anal., 13 (1982), pp. 459–483.
- [3] E. Cuesta, C. Lubich, and C. Palencia, Convolution quadrature time discretization of fractional diffusive-wave equations, Math. Comp., 75 (2006), pp. 673–696.
- [4] P. L. Duren, The Theory of H^p Spaces, Academic, New York, 1970.
- [5] G. GRIPENBERG, S-O. LONDEN, AND O. STAFFENS, Volterra Integral and Functional Equations, Cambridge University Press, Cambridge, 1990.
- [6] K. B. Hannsgen, Indirect Abelian theorems and a linear Volterra equation, Trans. Amer. Math. Soc., 142 (1969), pp. 539–555.
- [7] C. Lubich, Convolution quadrature and discretized operational calculus I, Numer. Math., 52 (1988), pp. 129–145.
- [8] C. LUBICH, I. H. SLOAN, AND V. THOMÉE, Nonsmooth data error estimates for approximations of an evolution equation with a positive-type memory term, Math. Comp., 65 (1996), pp. 1–
- [9] W. McLean and K. Mustapha, A second-order accurate numerical method for a fractional wave equation, Numer. Math., 105 (2007), pp. 481–510.
- [10] W. MCLEAN AND V. THOMÉE, Numerical solution of an evolution equation with a positive type memory term, J. Aust. Math. Soc., 35 (1993), pp. 23-70.
- [11] W. MCLEAN AND V. THOMEE, Numerical solution via Laplace transforms of a fractional order evolution equation, J. Integral Equations Appl., 22 (2010), pp. 57–94.

- [12] K. MUSTAPHA AND W. McLean, Discontinuous Galerkin method for an evolution equation with a memory term of positive type, Math. Comp., 78 (2009), pp. 1975–1995.
- [13] J. A. Nohel and D. F. Shea, Frequency domain methods for Volterra equations, Adv. Math., 22 (1976), pp. 278–304.
- [14] R. D. NOREN, Uniform L¹ behavior for the solution of a Volterra equation with a parameter, SIAM J. Math. Anal., 19 (1988), pp. 270–286.
- [15] J. PRÜSS, Evolutionary Integral Equations and Applications, Monogr. Math. 87, Birkhäuser, Basel, 1993.
- [16] M. RENARDY, W. J. HRUSA, AND J. A. NOHEL, Mathematical Problems in Viscoelasticity, Longman, London, 1987.
- 17 F. Riesz and B. Sz.-Nagy, Functional Analysis, Dover, New York, 1990.
- [18] W. Rudin, Real and Complex Analysis, McGraw-Hill, Boston, 1987.
- [19] D. F. SHEA AND S. WAINGER, Variants of the Wiener-Levy theorem with applications to stability problems for some Volterra integral equations, Amer. J. Math., 97 (1975), pp. 312–343.
- [20] D. V. Widder, The Laplace Transform, Princeton University Press, Princeton, NJ, 1946.
- [21] D. Xu, Uniform l¹ behavior for time discretization of a Volterra equation with completely monotonic kernel: I. Stability, IMA J. Numer. Anal., 22 (2002), pp. 133–151.
- [22] D. Xu, Uniform l¹ behavior for time discretization of a Volterra equation with completely monotonic kernel: II. Convergence, SIAM J. Numer. Anal., 46 (2008), pp. 231–259.
- [23] D. Xu, Stability of the difference type methods for linear Volterra equations in Hilbert space, Numer. Math., 109 (2008), pp. 571–595.
- [24] D. Xu, Uniform l¹ error bounds for the semidiscrete solution of a Volterra equation with completely monotonic convolution kernel, Comput. Math. Appl., 43 (2002), pp. 1303–1318.