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Measurement of Spin Density Matrix Elements in $\Lambda(1520)$ Photoproduction at 8.2-8.8 GeV

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Measurement of spin density matrix elements in $\Lambda(1520)$ photoproduction at 8.2–8.8 GeV

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We report on the measurement of spin density matrix elements of the $\Lambda(1520)$ in the photoproduction reaction $\gamma p \rightarrow \Lambda(1520)K^+$, via its subsequent decay to K^-p . The measurement was performed as part of the GlueX experimental program in Hall D at Jefferson Laboratory using a linearly polarized photon beam with $E_\gamma = 8.2\text{ GeV} - 8.8\text{ GeV}$. These are the first such measurements in this photon energy range. Results are presented in bins of momentum transfer squared, $-(t - t_0)$. We compare the results with a Reggeon exchange model and determine that natural exchange amplitudes are dominant in $\Lambda(1520)$ photoproduction.

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I. INTRODUCTION

The GlueX experiment is dedicated to expanding our knowledge of hadrons by measuring observables for a wide variety of states. The measurement presented here contributes to this effort by studying the photoproduction process of the $\Lambda(1520)$ hyperon ($J^P = 3/2^-$), specifically the measurement of spin density matrix elements (SDME). SDMEs parametrize the spin polarization of a produced state and are directly related to the underlying helicity amplitudes of the production process. As such they provide tests for scattering theory models which are needed in the search for new states in the hadron spectrum, especially in the search for small signals as are expected for exotic mesons. Schilling *et al.* [1] showed how SDMEs of vector mesons can be directly measured via the angular distribution of their decay products, and here we extend this technique for decays of spin- $3/2$ states. In addition to allowing us to measure the SDMEs of a strange baryon, the $\Lambda(1520)$ is experimentally attractive because it is a relatively isolated and narrow resonance with a width of $16\text{ MeV}/c^2$ [2].

The $\Lambda(1520)$ was discovered in 1962 using a K^- beam on a proton target [3], but since then only a few photoproduction measurements have been performed, with the majority of these at lower photon energies than the results reported in this paper. The only measurements performed in an energy range similar to that of GlueX are the differential cross sections from SLAC in 1971 [4] using an unpolarized photon beam. In 1980 the LAMP2 experiment extracted three independent SDMEs using an unpolarized photon beam with energy between 2.4 GeV and 4.8 GeV , in addition to differential cross sections [5]. The results indicated that the production does not proceed via simple K exchange. More recently, measurements at lower photon energies were published, mostly of cross sections [6–10].

Several attempts were made to describe the photoproduction of $\Lambda(1520)$ theoretically [11–14]. In general, the models used a Reggeon exchange approach to describe the t -channel production, which is expected to dominate beyond the s -channel resonance region. Since most of the available data cover a much lower energy range than that presented here, the models are not optimized for the GlueX energy range. Yu

and Kong [14], however, used the low energy results from LAMP2 [5] and high energy results from SLAC [4] to interpolate between available data and provide predictions for seven SDMEs in the GlueX energy range. In their model, they describe the production process in terms of K , K^* , and K_2^* exchanges, together with a proton pole in the s -channel and a contact term to preserve gauge invariance. They found that, especially at high energies, the K_2^* exchange is crucial to describe the data. Since the more recent cross section data from CLAS [9] disagrees with the LAMP2 [5] data by a factor of up to two at low energies (see Fig. 4 in Ref. [14]), Yu and Kong also made predictions from the same model based on the CLAS [9] and LEPS [6,7] data at lower energies. Both predictions will be used later to compare to our new data. Precise measurements of polarized SDMEs, such as those presented here, provide strong constraints on the production mechanisms used in models of $\Lambda(1520)$ photoproduction and will therefore help with our general understanding of photoproduction processes.

The structure of the paper is as follows. Section II introduces the SDMEs and gives the fit function used to extract them. Section III gives an overview of the experimental setup used for data taking. The event selection is presented in Sec. IV and Sec. V covers the methods used to extract the SDMEs from a sample of $\Lambda(1520)$ events. The results are discussed in Sec. VI.

II. FORMALISM

In order to study photoproduction of the $\Lambda(1520)$, we choose to reconstruct it in its decay to K^-p , which has a 22.5% branching fraction [2]. Therefore, we study the reaction $\gamma p \rightarrow K^+\Lambda(1520) \rightarrow K^+K^-p$.

We can learn about the production mechanism of the $\Lambda(1520)$ photoproduction by studying the spin transferred to it from the polarized photon. The spin density matrix ρ quantifies the spin polarization of the $\Lambda(1520)$ and parametrizes the angular distribution of its decay into K^-p . At high photon energies, t -channel exchange is expected to dominate this reaction, so it is convenient to study it in the Gottfried-Jackson (GJ), or t -channel helicity, system [15]. The coordinates are

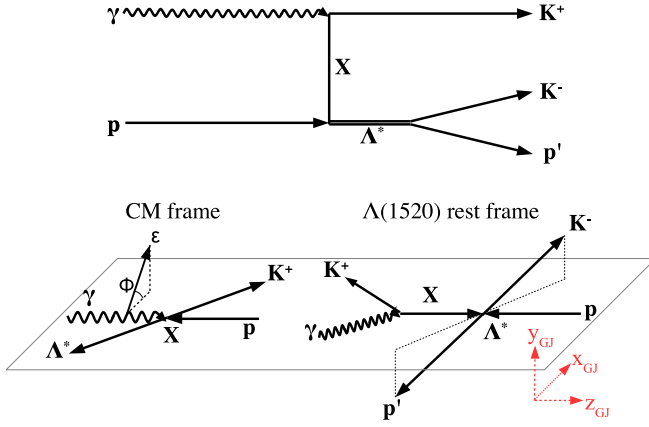


FIG. 1. The Gottfried-Jackson system used in this analysis. The diagram on top visualizes the t -channel production process expected to dominate at GlueX energies with X being the exchange particle. The Gottfried-Jackson system is defined in the rest frame of the $\Lambda(1520)$ [see Eq. (1)]. The polarization vector of the incoming beam photon is denoted by ϵ .

defined as

$$\hat{z} = \frac{-\vec{p}_p}{|-\vec{p}_p|}, \quad \hat{y} = \frac{\vec{p}_\gamma \times \vec{p}_{K^+}}{|\vec{p}_\gamma \times \vec{p}_{K^+}|}, \quad \hat{x} = \hat{y} \times \hat{z} \quad (1)$$

$$W(\theta, \phi, \Phi) = \frac{1}{2\pi} \frac{d\sigma}{dt} \frac{3}{4\pi} \left\{ \rho_{33}^0 \sin^2 \theta + \rho_{11}^0 \left(\frac{1}{3} + \cos^2 \theta \right) - \frac{2}{\sqrt{3}} \text{Re} \rho_{31}^0 \sin 2\theta \cos \phi - \frac{2}{\sqrt{3}} \text{Re} \rho_{3-1}^0 \sin^2 \theta \cos 2\phi \right. \\ \left. - P_\gamma \cos 2\Phi \left[\rho_{33}^1 \sin^2 \theta + \rho_{11}^1 \left(\frac{1}{3} + \cos^2 \theta \right) - \frac{2}{\sqrt{3}} \text{Re} \rho_{31}^1 \sin 2\theta \cos \phi - \frac{2}{\sqrt{3}} \text{Re} \rho_{3-1}^1 \sin^2 \theta \cos 2\phi \right] \right. \\ \left. - P_\gamma \sin 2\Phi \frac{2}{\sqrt{3}} \left[\text{Im} \rho_{31}^2 \sin 2\theta \sin \phi + \text{Im} \rho_{3-1}^2 \sin^2 \theta \sin 2\phi \right] \right\}. \quad (3)$$

In order to relate the spin of the particle to the production exchange mechanism, Schilling *et al.* showed that certain combinations of SDMEs can be expressed as linear combinations of purely natural or purely unnatural exchange amplitudes [1]. The naturality for a particle with spin-parity quantum number J^P is defined as $\eta = P(-1)^J$. As such, vector and tensor mesons [e.g., $K^*(892)$ and $K_2^*(1430)$] are natural exchanges ($\eta = +1$), and pseudoscalar and axial-vector mesons [e.g., $K(492)$ and $K_1(1270)$] are unnatural exchanges ($\eta = -1$). We denote production amplitudes for natural exchanges as N and for unnatural exchanges as U . Working in the reflectivity basis with helicities $\lambda_\gamma = \pm 1$, $\lambda_p = \pm 1/2$, and $\lambda_\Lambda = \pm 1/2, \pm 3/2$, and using the parity constraint results in four natural (N_σ) and four unnatural (U_σ) amplitudes, where $\sigma = \lambda_p - \lambda_\Lambda = \{-1, 0, 1, 2\}$:

$$\rho_{11}^0 + \rho_{11}^1 = \frac{2}{\mathcal{N}} (|N_0|^2 + |N_1|^2), \quad (4a)$$

$$\rho_{33}^0 + \rho_{33}^1 = \frac{2}{\mathcal{N}} (|N_{-1}|^2 + |N_2|^2), \quad (4b)$$

$$\text{Re}(\rho_{31}^0 + \rho_{31}^1) = \frac{2}{\mathcal{N}} \text{Re}(N_{-1}N_0^* - N_2N_1^*), \quad (4c)$$

with $\vec{p}_{p/\gamma/K^+}$ denoting the three-momentum of the target proton, incoming beam photon, and K^+ in the rest frame of the $\Lambda(1520)$. This is illustrated in Fig. 1. As the $\Lambda(1520)$ is a spin- $3/2$ particle, it has a 4×4 spin density matrix with 16 complex matrix elements. They are denoted by $\rho_{2\lambda_\Lambda, 2\lambda'_\Lambda}$, where λ_Λ denotes the $\Lambda(1520)$ helicity. Using a linearly polarized photon beam to produce the $\Lambda(1520)$ makes it possible to decompose the spin density matrix into

$$\rho = \rho^0 - P_\gamma \cos 2\Phi \rho^1 - P_\gamma \sin 2\Phi \rho^2, \quad (2)$$

where P_γ is the polarization of the photon beam and Φ is the angle between the photon polarization plane and the hadronic production plane, which is defined by the incoming γ and target proton and the outgoing K^+ and $\Lambda(1520)$ (see Fig. 1). Studying the decay $\Lambda(1520) \rightarrow K^- p$ means ten SDMEs, four unpolarized and six polarized, are accessible. To measure them, the distributions of θ and ϕ of the K^- in the GJ system are studied. These are given by Eq. (3) below [14]. This intensity distribution is normalized in such a way that integration over angles leads to the measured differential cross section $d\sigma/dt$, given the standard normalization $\rho_{33}^0 + \rho_{11}^0 = \frac{1}{2}$. There are thus nine independent SDMEs:

$$\text{Re}(\rho_{3-1}^0 + \rho_{3-1}^1) = \frac{2}{\mathcal{N}} \text{Re}(N_{-1}N_1^* + N_2N_0^*), \quad (4d)$$

$$\rho_{11}^0 - \rho_{11}^1 = \frac{2}{\mathcal{N}} (|U_0|^2 + |U_1|^2), \quad (4e)$$

$$\rho_{33}^0 - \rho_{33}^1 = \frac{2}{\mathcal{N}} (|U_{-1}|^2 + |U_2|^2), \quad (4f)$$

$$\text{Re}(\rho_{31}^0 - \rho_{31}^1) = \frac{2}{\mathcal{N}} \text{Re}(U_{-1}U_0^* - U_2U_1^*), \quad (4g)$$

$$\text{Re}(\rho_{3-1}^0 - \rho_{3-1}^1) = \frac{2}{\mathcal{N}} \text{Re}(U_{-1}U_1^* + U_2U_0^*). \quad (4h)$$

The normalization \mathcal{N} is given by

$$\mathcal{N} = 2(|N_{-1}|^2 + |N_0|^2 + |N_1|^2 + |N_2|^2 + |U_{-1}|^2 + |U_0|^2 + |U_1|^2 + |U_2|^2). \quad (4i)$$

These combinations can be used to study the naturality of exchanged particles in a t -channel exchange based on the extracted SDMEs. A full derivation of Eqs. (4a)–(4h) is given in Appendix A.

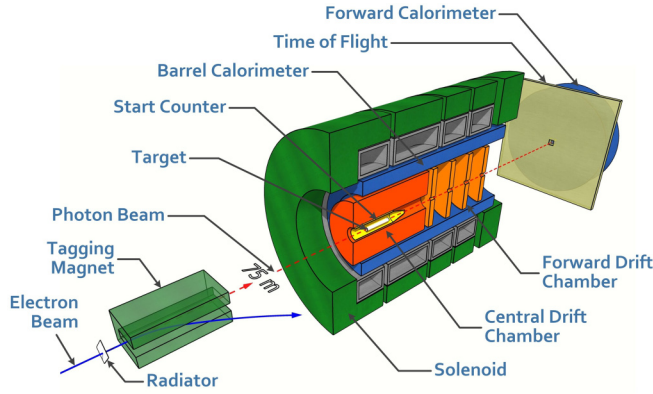


FIG. 2. Overview of the GlueX experiment and its important subdetector systems. Taken from Ref. [16].

III. GLUEX EXPERIMENT

The GlueX experiment is described in detail in Ref. [16]. A schematic overview is shown in Fig. 2. The GlueX spectrometer is located in Hall D at Jefferson Laboratory. To collect the data used in this paper, an 11.6 GeV electron beam provided by the Continuous Electron Beam Accelerator Facility (CEBAF) was used to produce a linearly polarized photon beam via the coherent bremsstrahlung technique on a thin diamond radiator. The orientation of the beam polarization plane was controlled by adjusting the orientation of the diamond using a goniometer. During data taking, four different pairwise orthogonal diamond settings were used in turn to control systematic effects. The beam polarization had its maximum in the coherent peak, whose position was also controlled through diamond orientation. Figure 3 shows the degree of polarization, as measured with a triplet polarimeter [17], for the four different diamond orientations. The measurement of the polarization carries a systematic uncertainty of 1.5% [17]. Together with a 3% statistical uncertainty, this results in an overall uncertainty on the degree of linear polarization of $\pm 3.5\%$. Only events with a photon beam energy in the range from

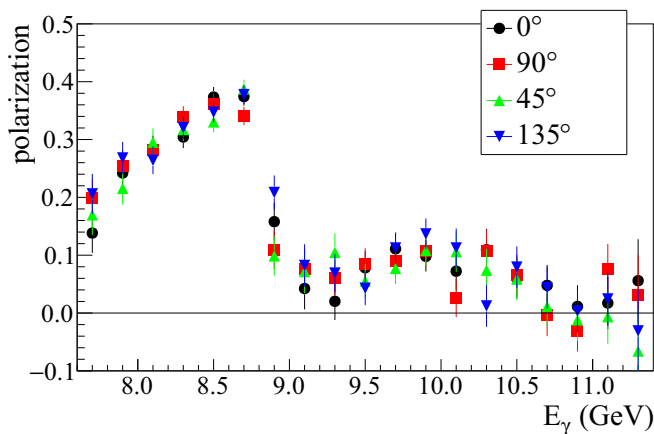


FIG. 3. The photon beam polarization is shown for the four separate diamond settings. The measurement carries a systematic uncertainty of 1.5% [17].

$E_\gamma = 8.2$ GeV to 8.8 GeV, where the polarization and also the flux were highest, were analyzed. For each diamond setting, the average polarization in this range was determined and used for further analysis. In addition, for about 15% of the data, an aluminum radiator was used to generate an unpolarized photon beam. Measurements using each of the beam settings (four polarized, one unpolarized) were distributed evenly across the beamtime to minimize systematic effects such as small drifts in detector acceptance or efficiency.

The electrons scattering from the radiator were deflected by a dipole magnet onto the tagger focal plane, where an array of scintillation detectors measured their position, and hence momentum, allowing the energy of associated bremsstrahlung photons to be determined. The collimated photon beam was incident on the liquid hydrogen target, which was enclosed by the start counter (SC), a scintillation detector that provides a reference time for each event. Surrounding this were the central drift chamber (CDC) for tracking of charged particles, and the barrel calorimeter (BCAL), which also enclosed the forward drift chamber (FDC), all in a 2 T magnetic field. The tracking detectors had a momentum resolution of $\sigma_p/p \approx 1-5\%$. A time-of-flight detector, with the main purpose of particle identification, and the forward calorimeter (FCAL), were placed in the forward direction. The excellent timing resolution of the BCAL of 150 ps at 1 GeV meant that it was possible to use it for a time-of-flight measurement to help with particle identification (PID).

The data used in this measurement were collected in spring 2017 and correspond to an integrated luminosity of about 21.8 pb^{-1} . The main readout trigger required a minimum energy deposition in either the BCAL or a combination of BCAL and FCAL.

To model the detector acceptance and reconstruction efficiency, a standardized GEANT4-based [18] GlueX detector simulation, *hdgeant4*, was used [16].

IV. EVENT SELECTION

In order to select the reaction $\gamma p \rightarrow K^+ \Lambda(1520) \rightarrow K^+ K^- p$, events with at least two positively charged particles and one negatively charged final-state particle were analyzed. Up to three additional charged tracks were allowed to be detected in an event, to make sure that good events were not erroneously rejected because of spurious tracks in the detector. Each combination of two positive tracks and one negative track was analyzed as $K^+ K^- p$. For particle identification, time-of-flight requirements were placed for each track, using the detector with the best available timing information. Furthermore, the energy loss dE/dx of the proton in the CDC was used for PID. A kinematic fit was carried out with the fit hypothesis $\gamma p \rightarrow K^+ K^- p$, which included vertex and four-momentum constraints. Events with a kinematic fit confidence level of $\text{CL} < 10^{-6}$ were rejected. Also, it was required that the particle tracks originated from within the target cell. To restrict events to the $\Lambda(1520)$ signal region, only those with a pK^- invariant mass between $1.46 \text{ GeV}/c^2$ and $1.58 \text{ GeV}/c^2$ were analyzed (see Fig. 4).

The electron beam, and hence the photon beam, had a bunch structure with a timing separation of 4.008 ns. Each

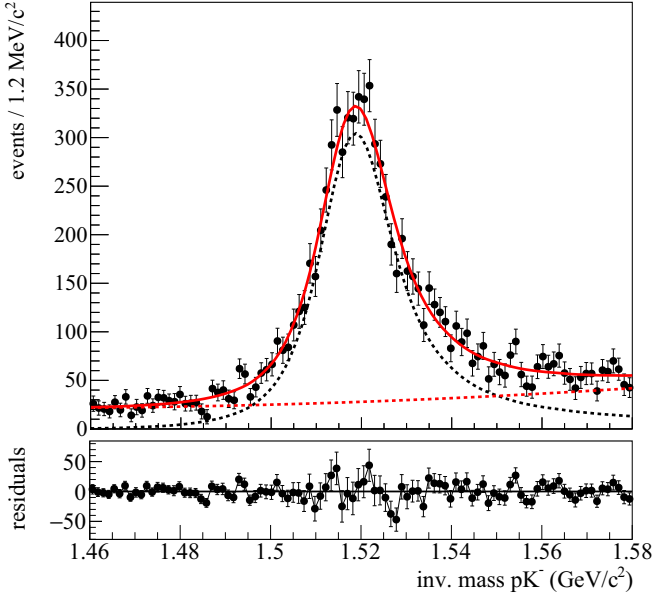


FIG. 4. Example sPlot fit [19] for one bin between $-(t - t_0) = 0.3 \text{ GeV}^2/c^2 - 0.5 \text{ GeV}^2/c^2$. Top: The black and red dashed lines show the fit components of signal and background respectively, the red solid line shows the resulting fit to the data (black points). Bottom: Residuals of the total fit to the data.

bunch resulted in multiple hits on the tagger, of which only one belonged to the beam photon that triggered the event. This beam photon was determined via a coincidence between the hadronic event time and the bunch time. To remove the background from photons within the same bunch, a statistical sideband subtraction was performed. For this, tagger hits that were recorded close in time before and after the bunch in coincidence with the hadronic event were analyzed. These events were given a negative weight proportional to the relative size of coincidence peak and sideband regions, which were defined based on the time difference between the tagger hit and the beam bunch.

The sPlot technique [19], which was successfully used in other experiments extracting polarization observables [20,21], was used to subtract the remaining background under the $\Lambda(1520)$ signal peak by determining event-by-event sWeights, which, when applied to the data, resulted in the signal distribution used for further analysis. This was done in eight bins of four-momentum transfer $-(t - t_0)$, where t is the Mandelstam variable that describes the transfer of four-momentum from the beam photon to the target proton. Its kinematic limit is given by t_0 . To create a signal shape, events were generated according to a relativistic Breit-Wigner distribution and then simulated with *hdgeant4*. Additional parameters were added to the signal function to allow flexibility in accounting for small differences between data and simulation. A second-degree Chebyshev polynomial was used to parametrize the background under the $\Lambda(1520)$. The fit was performed within the *brufit* framework [22], which uses ROOFIT [23]. An example fit is shown in Fig. 4. The dashed black and red lines show the signal and background contributions, respectively. The solid red line is the total fit to the data (black points).

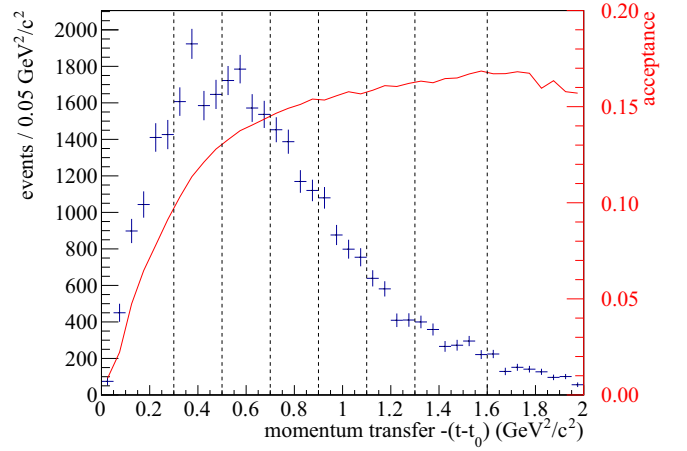


FIG. 5. Distribution of momentum transfer $-(t - t_0)$ for signal weighted events (blue data points) and acceptance (red line). The black dotted lines indicate the chosen bin limits.

In addition, the fit residuals, which are used to assess the fit quality, are shown. The plots show that the chosen distributions describe the data very well. To test a potential impact of the background model on the results, three additional variations of Chebyshev polynomials were tested. No significant systematic effect was observed.

After applying the sPlot background subtraction, about 32,200 events remained for the extraction of SDMEs. Their $-(t - t_0)$ distribution is shown in Fig. 5. The dashed black lines indicate the bin limits used in this analysis. The solid red line represents the acceptance as determined from simulations.

V. SDME PARAMETER ESTIMATION

To estimate the nine independent SDME parameters, the Markov chain Monte Carlo (MCMC) technique was used [24]. Instead of minimizing a χ^2 or negative log-likelihood, as is often done, this method explores the possible parameter space numerically using the Metropolis-Hastings (MH) algorithm [25,26]. For this purpose, a likelihood function was written as

$$\ln \mathcal{L} = s_w \left(\sum_{i=1}^N s w_i \ln \mathcal{I} - \int d\Omega \mathcal{I} \eta(\Omega) \right) \quad (5)$$

with $\mathcal{I} = W(\theta, \phi, \Phi)$ [Eq. (3)] being the intensity function. The sWeights are notated by $s w_i$, and

$$s_w = \frac{\sum_{i=1}^N s w_i}{\sum_{i=1}^N s w_i^2} \quad (6)$$

is a constant factor accounting for the effect of the weights on the statistical uncertainty. While the sum in the likelihood ran over all N events in the dataset, the integral was evaluated as a sum over simulated data that were generated flat in phase space and processed through *hdgeant4*. This accounted for detector acceptance effects denoted as $\eta(\Omega)$ in the likelihood function.

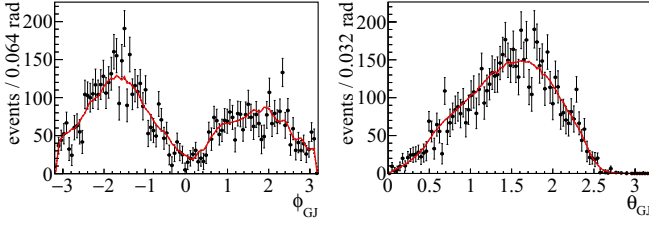


FIG. 6. Projections of data (black points) and simulations containing acceptance effects which were weighted with fit results (red lines) for $-(t - t_0) = 0.3 \text{ GeV}^2/c^2 - 0.5 \text{ GeV}^2/c^2$. Note, as the red curves are based on simulated data, they are not expected to be perfectly smooth.

For this analysis, the MH implementation of ROOSTATS [27] was used. As prior, a uniform distribution of SDME values with range $[-1, 1]$ was assumed, reflecting the fact that SDMEs are confined to this region. New steps in the Markov chain were proposed by a *sequential proposal* function which randomly changed one of the nine SDME parameters at a time and proposed its next value based on a Gaussian distribution centered around the current value with a width tuned to achieve an acceptance rate of about 10–20%. Final parameters are reported as the means of the posterior distributions, with uncertainties given by the widths.

In order to visually assess the quality of the extracted SDME parameters, simulations that were produced flat in phase space and processed through *hdgeant4* to incorporate detector inefficiencies were reweighted with the resulting intensity function. A comparison of this weighted simulation to data, for one example bin and the two variables θ_{GJ} and ϕ_{GJ} , is shown in Fig. 6.

A. Validation

Extensive studies on simulated data have been performed to validate that the presented approach to extract the SDME parameters gave on average the correct results and uncertainties for the estimated parameters. Each study included simulations of 400 statistically independent samples of signal and background with preselected, known SDMEs. The simulations were processed through the complete *hdgeant4* simulation and treated as real data from that point onwards, including the full event selection and sPlot background subtraction. For each sample, the SDMEs were extracted and compared to the parameters chosen for generation. Observed differences were quantified in a systematic uncertainty (see next section). Details on all studies can be found in Ref. [28].

B. Systematic uncertainties

As described earlier, the relative systematic uncertainty on the beam polarization was determined to be $\pm 3.5\%$. This only affected the polarized $\rho^{1,2}$ SDMEs and the uncertainties are fully correlated across the full $-(t - t_0)$ range. The absolute systematic uncertainty of the extraction method of SDME parameters, obtained from the validation studies showed only very little correlation across the $-(t - t_0)$ range and was

TABLE I. Summary of systematic uncertainties. The uncertainties for extraction method and simulation model are absolute numbers while the uncertainty on the degree of polarization is a relative scaling uncertainty.

Source	Uncertainty ρ^0	Uncertainty $\rho^{1,2}$
Extraction method	0.02	0.01
Simulation model	0.01	0.007
Degree of polarization	–	3.5%

± 0.02 for unpolarized and ± 0.01 for polarized SDMEs. In addition to these and the assumed background distribution in the sPlot fit, other aspects of the analysis were tested carefully for systematic effects. Twenty-six different variations in the event selection, including different limits for the pK^- invariant mass range, the kinematic fit confidence level, vertex position, and timing of the particles in the BCAL and TOF were considered. None of them showed significant systematic effects on the results. The full list of tested variations can be found in Ref. [28].

In order to explore systematic uncertainties, due to inaccuracies in the simulated model of the angular dependence of the tracking efficiency, we used the results of a study similar to that described in Sec. 15.1 of Ref. [16]. For each of the three tracks in the event we obtained the ratio of efficiency in data to efficiency in simulation for the particular region in the two-dimensional momentum- θ plane. For tracks beyond the region of phase space covered in Ref. [16], we used their largest measured ratio. We then reweighted the accepted simulated data by the product of these weights for each of the three tracks and repeated the SDME analysis. We observed that the absolute central values of the SDMEs changed by no more than ± 0.01 for the unpolarized ρ^0 and ± 0.007 for the polarized $\rho^{1,2}$, with little correlation across the $-(t - t_0)$ range. We therefore used this as our estimate for the systematic uncertainty due to inaccuracies in the simulated model and combine it with the other systematic uncertainties in quadrature.

A summary of the relevant systematic uncertainties is given in Table I. The total systematic uncertainty for each individual bin is given in Table II together with the results.

VI. RESULTS

Results are shown in Fig. 7. The vertical error bars show the statistical uncertainty, the blue shaded boxes the scaling uncertainty from the polarization, and the black box the remaining systematic uncertainties combined in quadrature. The horizontal error bars show the root mean square (rms) widths within the $-(t - t_0)$ bins. Also shown in the figure are predictions made by Yu and Kong (private communication based on Ref. [14]) for seven of the nine extracted SDMEs. The blue solid lines show the predictions based on data from CLAS [9] and LEPS [6,7], and the red dashed lines show predictions based on data from LAMP2 [5] and SLAC [4]. These predictions are based on previous data at much lower or higher photon beam energies, and until now there have been no data for polarized SDMEs. It is clear that these new GlueX

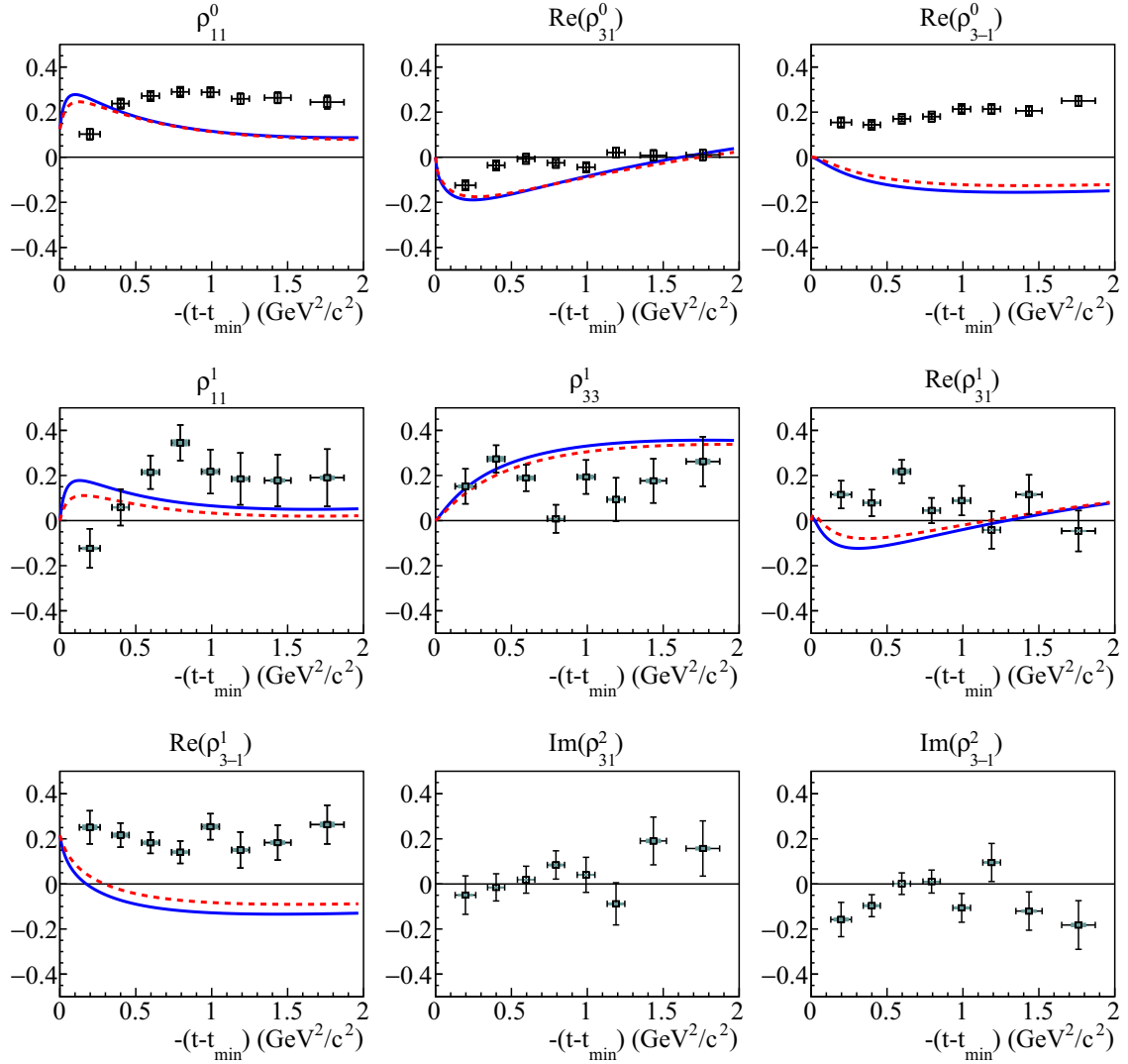


FIG. 7. Spin density matrix elements and predictions by Yu and Kong (based on Ref. [14]), using parameters based on data from CLAS [9] and LEPS [6,7] (blue solid) and using parameters based on data from LAMP2 [5] and SLAC [4] (red dashed). The vertical error bars show the statistical uncertainty, the blue shaded boxes the scaling uncertainty from the polarization, and the black boxes the the remaining systematic uncertainties combined in quadrature. The horizontal error bars show the rms widths within the $-(t - t_0)$ bins.

data will place stringent new constraints on the model. To interpret the extracted SDMEs in terms of the contributing exchange mechanism, the combinations from Eqs. (4a)–(4h) were formed and are shown in Fig. 8. Over most of the $-(t - t_0)$ range, the results indicate natural exchanges are dominant. Only in the lowest bin, there seems to be a small contribution from unnatural exchanges. Although the observed dominance of natural exchange does not allow us to specify a particular exchange particle, we note that Yu and Kong predicted a dominant natural exchange at high energies, via a K_2^* . This can be seen by looking at the same combinations for their predictions. The unnatural contributions are 0 over most of the $-(t - t_0)$ range. While their model does not agree well with the present SDME measurements, the expectation of a large natural contribution to the exchange is supported by the data.

VII. CONCLUSION

Nine independent spin density matrix elements for the reaction $\gamma p \rightarrow K^+ \Lambda(1520) \rightarrow K^+ K^- p$ have been measured for photon beam energies between 8.2 GeV to 8.8 GeV. For the $\Lambda(1520)$ this represents the first measurement of polarized SDMEs and the first measurement of unpolarized SDMEs at these energies. Our measurements allow conclusions about the production mechanisms by studying combinations of SDMEs that represent purely natural or unnatural exchanges in a t -channel exchange. It was found that the photoproduction of $\Lambda(1520)$ at these energies is dominated by natural exchange amplitudes over most of the $-(t - t_0)$ range, in qualitative agreement with the only available model prediction. However, the quantitative agreement with the model was poor. The presented results will hopefully motivate more work on the photoproduction of $\Lambda(1520)$ which should lead to better

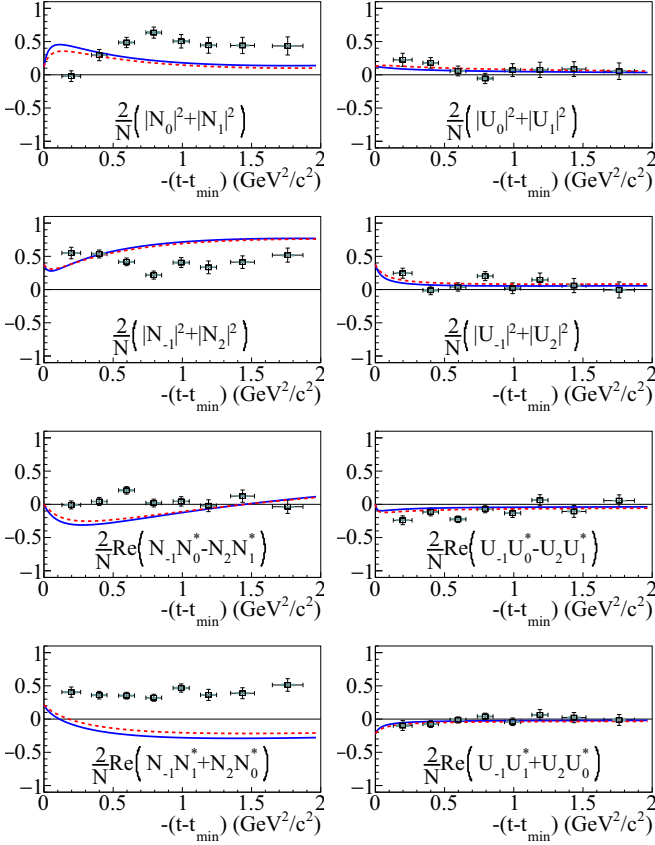


FIG. 8. Combinations of SDMEs corresponding to natural (left column) and unnatural (right column) exchanges in the photoproduction of $\Lambda(1520)$. Also shown are the corresponding combinations of predictions by Yu and Kong (based on Ref. [14]), using parameters based on data from CLAS [9] and LEPS [6,7] (blue solid) and using parameters based on data from LAMP2 [5] and SLAC [4] (red dashed). The vertical error bars show the statistical uncertainty, the blue shaded boxes the scaling uncertainty from the polarization, and the black boxes the remaining systematic uncertainties combined in quadrature. The horizontal error bars show the rms widths within the $-(t-t_0)$ bins.

agreement between data and models in the future. To further the understanding of this reaction, precise measurements of differential cross sections are desirable. GlueX is ideally placed to perform these over a wide range of energies.

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APPENDIX A: DERIVATION OF NATURAL AND UNNATURAL AMPLITUDES

We provide a brief derivation of Eqs. (4a)–(4h), which are used to interpret the SDMEs in terms of the naturality $\eta = P(-1)^J$ of the exchanged particle.

Following Schilling [1], we denote the production amplitude as $T_{\lambda_\gamma \lambda_p \lambda_\Lambda}$ with helicities $\lambda_\gamma = \pm 1$, $\lambda_p = \pm 1/2$, and $\lambda_\Lambda = \pm 1/2, \pm 3/2$. Taking parity $T_{-\lambda_\gamma -\lambda_p -\lambda_\Lambda} = (-1)^{\lambda_p - \lambda_\Lambda} T_{\lambda_\gamma \lambda_p \lambda_\Lambda}$ into account, this leaves us with eight independent amplitudes. We can split each amplitude into a positive (N , $\epsilon = +1$) and a negative (U , $\epsilon = -1$) component and write

$$T = T^{(+)} + T^{(-)} \quad (\text{A1})$$

with the amplitudes given in reflectivity basis by

$$T_{\lambda_p \lambda_\Lambda}^{(\epsilon)} = \frac{1}{2}(T_{1\lambda_p \lambda_\Lambda} + \epsilon T_{-1\lambda_p \lambda_\Lambda}). \quad (\text{A2})$$

We can use these amplitudes to express the SDMEs as

$$\rho_{\lambda_\Lambda \lambda'_\Lambda}^0 = \frac{1}{\mathcal{N}} \sum_{\lambda_p} T_{\lambda_p \lambda_\Lambda}^{(+)} T_{\lambda_p \lambda'_\Lambda}^{(+)*} + T_{\lambda_p \lambda_\Lambda}^{(-)} T_{\lambda_p \lambda'_\Lambda}^{(-)*}, \quad (\text{A3})$$

$$\rho_{\lambda_\Lambda \lambda'_\Lambda}^1 = \frac{1}{\mathcal{N}} \sum_{\lambda_p} T_{\lambda_p \lambda_\Lambda}^{(+)} T_{\lambda_p \lambda'_\Lambda}^{(+)*} - T_{\lambda_p \lambda_\Lambda}^{(-)} T_{\lambda_p \lambda'_\Lambda}^{(-)*}, \quad (\text{A4})$$

$$\rho_{\lambda_\Lambda \lambda'_\Lambda}^2 = \frac{i}{\mathcal{N}} \sum_{\lambda_p} T_{\lambda_p \lambda_\Lambda}^{(+)} T_{\lambda_p \lambda'_\Lambda}^{(-)*} - T_{\lambda_p \lambda_\Lambda}^{(-)} T_{\lambda_p \lambda'_\Lambda}^{(+)*}, \quad (\text{A5})$$

and write the following combinations:

$$\rho_{\lambda_\Lambda \lambda'_\Lambda}^0 + \rho_{\lambda_\Lambda \lambda'_\Lambda}^1 = \frac{2}{\mathcal{N}} \sum_{\lambda_p} T_{\lambda_p \lambda_\Lambda}^{(+)} T_{\lambda_p \lambda'_\Lambda}^{(+)*}, \quad (\text{A6})$$

$$\rho_{\lambda_\Lambda \lambda'_\Lambda}^0 - \rho_{\lambda_\Lambda \lambda'_\Lambda}^1 = \frac{2}{\mathcal{N}} \sum_{\lambda_p} T_{\lambda_p \lambda_\Lambda}^{(-)} T_{\lambda_p \lambda'_\Lambda}^{(-)*}. \quad (\text{A7})$$

These combinations separate the amplitudes with positive and negative reflectivity. We can further write the amplitudes as N_σ (natural) and U_σ (unnatural) with $\sigma = \lambda_p - \lambda_\Lambda = \{-1, 0, 1, 2\}$ and write the reflectivity amplitudes in terms of the exchange naturality

$$N_{-1} = T_{\frac{1}{2}\frac{3}{2}}^{(+)}, \quad N_0 = T_{\frac{1}{2}\frac{1}{2}}^{(+)}, \quad N_1 = T_{\frac{1}{2}\frac{-1}{2}}^{(+)}, \quad N_2 = T_{\frac{1}{2}\frac{-3}{2}}^{(+)}, \quad (\text{A8})$$

$$U_{-1} = T_{\frac{1}{2}\frac{3}{2}}^{(-)}, \quad U_0 = T_{\frac{1}{2}\frac{1}{2}}^{(-)}, \quad U_1 = T_{\frac{1}{2}\frac{-1}{2}}^{(-)}, \quad U_2 = T_{\frac{1}{2}\frac{-3}{2}}^{(-)}. \quad (\text{A9})$$

Using Eqs. (A8) and (A9), with Eqs. (A6) and (A7), leads directly to Eqs. (4a)–(4h) with normalization given by Eq. (4i).

TABLE II. Numerical results for all presented SDMEs, natural and unnatural combinations, and covariances between ρ_{11}^1 and ρ_{33}^1 . The first uncertainty is statistical, the second systematic.

$-(t-t_0)$ in GeV ² /c ²	ρ_{11}^0	ρ_{31}^0	ρ_{3-1}^0	ρ_{11}^1	ρ_{33}^1	ρ_{31}^1	ρ_{3-1}^1	ρ_{31}^2	ρ_{3-1}^2
0.197 ± 0.069	0.102 ± 0.025 ± 0.022	-0.125 ± 0.016 ± 0.022	0.154 ± 0.018 ± 0.022	-0.123 ± 0.087 ± 0.013	0.152 ± 0.078 ± 0.013	0.115 ± 0.061 ± 0.013	0.251 ± 0.074 ± 0.015	-0.049 ± 0.085 ± 0.012	-0.158 ± 0.076 ± 0.013
0.400 ± 0.056	0.238 ± 0.016 ± 0.022	-0.036 ± 0.015 ± 0.022	0.144 ± 0.013 ± 0.022	0.059 ± 0.081 ± 0.012	0.273 ± 0.061 ± 0.016	0.078 ± 0.059 ± 0.013	0.217 ± 0.054 ± 0.014	-0.015 ± 0.060 ± 0.012	-0.096 ± 0.049 ± 0.013
0.597 ± 0.057	0.272 ± 0.015 ± 0.022	-0.007 ± 0.013 ± 0.022	0.170 ± 0.012 ± 0.022	0.214 ± 0.074 ± 0.014	0.189 ± 0.058 ± 0.014	0.218 ± 0.051 ± 0.014	0.183 ± 0.047 ± 0.014	0.019 ± 0.060 ± 0.012	0.001 ± 0.048 ± 0.012
0.793 ± 0.057	0.290 ± 0.016 ± 0.022	-0.025 ± 0.013 ± 0.022	0.180 ± 0.012 ± 0.022	0.345 ± 0.079 ± 0.017	0.008 ± 0.062 ± 0.012	0.045 ± 0.056 ± 0.012	0.141 ± 0.050 ± 0.013	0.084 ± 0.063 ± 0.013	0.011 ± 0.051 ± 0.012
0.992 ± 0.058	0.288 ± 0.019 ± 0.022	-0.044 ± 0.017 ± 0.022	0.213 ± 0.014 ± 0.022	0.217 ± 0.098 ± 0.014	0.193 ± 0.075 ± 0.014	0.088 ± 0.065 ± 0.013	0.254 ± 0.058 ± 0.015	0.040 ± 0.078 ± 0.012	-0.106 ± 0.064 ± 0.013
1.189 ± 0.058	0.259 ± 0.024 ± 0.022	0.021 ± 0.021 ± 0.022	0.214 ± 0.018 ± 0.022	0.185 ± 0.115 ± 0.014	0.094 ± 0.096 ± 0.014	-0.042 ± 0.083 ± 0.012	0.150 ± 0.080 ± 0.013	-0.088 ± 0.094 ± 0.013	0.095 ± 0.084 ± 0.013
1.435 ± 0.086	0.264 ± 0.025 ± 0.022	0.008 ± 0.024 ± 0.022	0.206 ± 0.019 ± 0.022	0.178 ± 0.114 ± 0.014	0.176 ± 0.098 ± 0.014	0.115 ± 0.087 ± 0.013	0.183 ± 0.077 ± 0.014	0.191 ± 0.106 ± 0.014	-0.120 ± 0.085 ± 0.013
1.761 ± 0.111	0.244 ± 0.030 ± 0.022	0.010 ± 0.024 ± 0.022	0.250 ± 0.021 ± 0.022	0.190 ± 0.127 ± 0.014	0.262 ± 0.110 ± 0.015	-0.046 ± 0.091 ± 0.012	0.263 ± 0.086 ± 0.015	0.157 ± 0.122 ± 0.013	-0.182 ± 0.108 ± 0.014
$-(t-t_0)$ in GeV ² /c ²	$\sigma_{\rho_{11}^1 \rho_{33}^1}$	$\frac{2}{\mathcal{N}}(M_0 ^2 + M_1 ^2)$	$\frac{2}{\mathcal{N}}(N_{-1} ^2 + N_2 ^2)$	$\frac{2}{\mathcal{N}}\text{Re}(N_{-1}N_0^* - N_2N_1^*)$	$\frac{2}{\mathcal{N}}\text{Re}(N_{-1}N_1^* + N_2N_0^*)$	$\frac{2}{\mathcal{N}}(U_0 ^2 + U_1 ^2)$	$\frac{2}{\mathcal{N}}(U_{-1} ^2 + U_2 ^2)$	$\frac{2}{\mathcal{N}}\text{Re}(U_{-1}U_0^* - U_2U_1^*)$	$\frac{2}{\mathcal{N}}\text{Re}(U_{-1}U_1^* + U_2U_0^*)$
0.197 ± 0.069	-0.00438	-0.021 ± 0.081 ± 0.035	0.550 ± 0.087 ± 0.035	-0.010 ± 0.060 ± 0.035	0.405 ± 0.077 ± 0.036	0.225 ± 0.099 ± 0.035	0.246 ± 0.076 ± 0.035	-0.240 ± 0.066 ± 0.035	-0.097 ± 0.075 ± 0.036
0.400 ± 0.056	-0.00368	0.296 ± 0.085 ± 0.035	0.536 ± 0.060 ± 0.036	0.043 ± 0.061 ± 0.035	0.360 ± 0.057 ± 0.035	0.179 ± 0.080 ± 0.035	-0.011 ± 0.066 ± 0.036	-0.114 ± 0.059 ± 0.035	-0.073 ± 0.053 ± 0.035
0.597 ± 0.057	-0.00321	0.486 ± 0.077 ± 0.035	0.417 ± 0.058 ± 0.035	0.210 ± 0.057 ± 0.035	0.353 ± 0.051 ± 0.035	0.058 ± 0.074 ± 0.035	0.039 ± 0.062 ± 0.035	-0.225 ± 0.050 ± 0.035	-0.013 ± 0.046 ± 0.035
0.793 ± 0.057	-0.00363	0.635 ± 0.084 ± 0.037	0.218 ± 0.061 ± 0.034	0.019 ± 0.059 ± 0.035	0.321 ± 0.053 ± 0.035	-0.055 ± 0.077 ± 0.037	0.202 ± 0.068 ± 0.034	-0.070 ± 0.057 ± 0.035	0.039 ± 0.050 ± 0.035
0.992 ± 0.058	-0.00544	0.505 ± 0.102 ± 0.035	0.406 ± 0.075 ± 0.035	0.044 ± 0.070 ± 0.035	0.467 ± 0.062 ± 0.036	0.070 ± 0.097 ± 0.035	0.019 ± 0.080 ± 0.035	-0.133 ± 0.065 ± 0.035	-0.041 ± 0.057 ± 0.036
1.189 ± 0.058	-0.00772	0.444 ± 0.120 ± 0.035	0.334 ± 0.096 ± 0.035	-0.021 ± 0.089 ± 0.035	0.364 ± 0.084 ± 0.035	0.074 ± 0.115 ± 0.035	0.147 ± 0.102 ± 0.035	0.062 ± 0.083 ± 0.035	0.063 ± 0.080 ± 0.035
1.435 ± 0.086	-0.00794	0.441 ± 0.123 ± 0.035	0.412 ± 0.095 ± 0.035	0.124 ± 0.092 ± 0.035	0.389 ± 0.082 ± 0.035	0.086 ± 0.111 ± 0.035	0.060 ± 0.107 ± 0.035	-0.107 ± 0.089 ± 0.035	0.022 ± 0.077 ± 0.035
1.761 ± 0.111	-0.00941	0.434 ± 0.136 ± 0.035	0.518 ± 0.106 ± 0.036	-0.036 ± 0.100 ± 0.035	0.513 ± 0.095 ± 0.036	0.054 ± 0.124 ± 0.035	-0.006 ± 0.121 ± 0.036	0.056 ± 0.087 ± 0.035	-0.013 ± 0.082 ± 0.036

APPENDIX B: NUMERICAL RESULTS

All numerical results for the SDMEs and their statistical and systematic uncertainties, together with the natural and unnatural combinations are listed in Table II. In general, the correlations in the statistical uncertainties are small, except for

ρ_{11}^1 and ρ_{33}^1 , whose covariances need to be taken into account when using the data further, and are listed as well.

Subsets of the Markov chains used for the parameter estimation are available as Supplemental Material [29].

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